

Millimeter interferometers

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6th IRAM Millimeter Interferometry School 6–10 October 2008



Millimeter interferometers Outline

- Goal: the van Cittert–Zernike theorem
- The ideal interferometer
 - \hookrightarrow geometrical delay, source size, bandwidth
- The real interferometer
 - \hookrightarrow heterodyne receivers, delay correction, correlators
- Aperture synthesis

 $\hookrightarrow uv$ plane, field of view, transfer function

• Sensitivity



Goal van Cittert–Zernike theorem

- \bullet Wiener-Kichnine theorem
 - autocorrelation of S(t) = FT(spectra) $S(t_1) S(t_2) = AS(\tau) \rightleftharpoons S(\nu)$
 - -implementation: FT spectrometers
- van Cittert–Zernike theorem
 - -spatial autocorrelation of S(x) = FT(brightness) $S(x_1) S(x_2) = AS(u) \rightleftharpoons S(\alpha)$
 - implementation: interferometry



Goal van Cittert–Zernike theorem

Implementing the van Cittert–Zernike theorem

- 1. Build a device that measures the spatial autocorrelation of the incoming signal
- 2. Do it for all possible scales
- 3. Take the FT and get an image of the brightness distribution



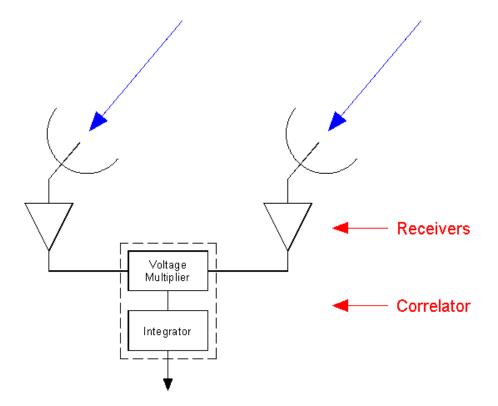
Goal Interferometry

Implementing the van Cittert–Zernike theorem

- 1. Build a device that measures the spatial autocorrelation of the incoming signal \longrightarrow 2-elements interferometer
- 2. Do it for all possible scales $\longrightarrow \mathbf{N}$ antennas
- 3. Take the FT and get an image of the brightness distribution \longrightarrow **software**



The ideal interferometer Sketch





- The heterodyne <u>receiver</u> measures the incoming electric field $E \cos(2\pi\nu t)$
- The <u>correlator</u> is a <u>multiplier</u> followed by a <u>time integrator</u>:

 $r = \langle E_1 \cos(2\pi\nu t) | E_2 \cos(2\pi\nu t) \rangle = E_1 E_2$

• We have measured the spatial correlation of the signal!



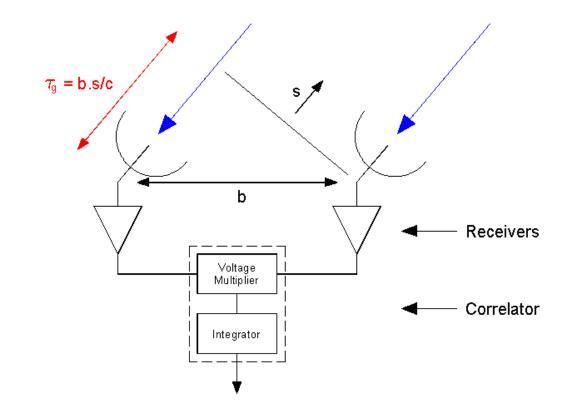
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 $r = \langle E_1 \cos(2\pi\nu t) | E_2 \cos(2\pi\nu t) \rangle = E_1 E_2$

- We have measured the spatial correlation of the signal!
- But we have forgotten the geometrical delay



The ideal interferometer Sketch





- There is a **geometrical delay** τ_g between the two antennas \longrightarrow **more complex** experiment than the Young's hole
- Correlator output:

 $\begin{aligned} r &= < E_1 \cos(2\pi\nu t) \ E_2 \cos(2\pi\nu t) > = E_1 \ E_2 \\ r &= < E_1 \cos(2\pi\nu (t-\tau_g)) \ E_2 \cos(2\pi\nu t) > \\ &= E_1 E_2 \cos(2\pi\nu \tau_g) \end{aligned}$



- Correlator output: $r = E_1 E_2 \cos(2\pi\nu\tau_g)$
- τ_g varies slowly with time (Earth rotation) \longrightarrow fringes
- Natural fringe rate:

 \sim

$$\tau_g = \frac{\mathbf{b.s}}{c} \qquad \nu \ \frac{d\tau_g}{dt} \simeq \Omega_{earth} \ \frac{\mathbf{b}\nu}{c}$$

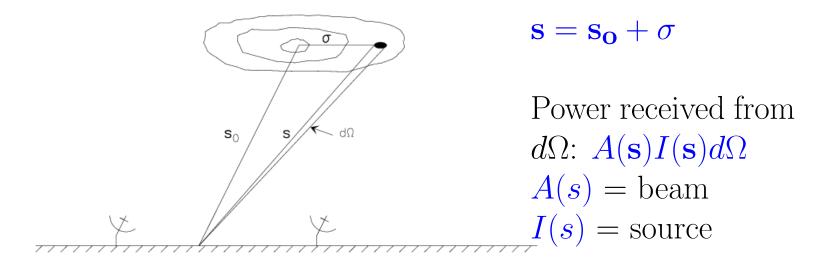
$$\nu \ 50 \text{ Hz for } b = 800 \text{ m and } \nu = 250 \text{ GHz}$$



- Correlator output: $r = E_1 E_2 \cos(2\pi\nu\tau_g)$
- τ_g varies slowly with time (Earth rotation) \longrightarrow fringes
- τ_g is **known** from the antenna position, source direction, time \longrightarrow could be corrected
- Problems: the source is **not a point source** $(\int \tau_g ?)$ the signal is **not monochromatic** $(\int \nu ?)$



The ideal interferometer Source size



Correlator output: $r = E_1 E_2 \cos(2\pi\nu\tau_g)$ $r = A(\mathbf{s})I(\mathbf{s})d\Omega\cos(2\pi\nu\tau_g(\mathbf{s}))$



The ideal interferometer Source size

• Correlator output (integrated):

$$R = \int_{Sky} A(\mathbf{s})I(\mathbf{s})\cos(2\pi\nu\mathbf{b}.\mathbf{s}/c) d\Omega$$
$$= |V|\cos(2\pi\nu\tau_g - \varphi_V)$$

• Complex visibility:

$$V = |V|e^{i\varphi_{\rm V}} = \int_{Sky} A(\sigma)I(\sigma)e^{-2i\pi\nu\mathbf{b}.\sigma/c}d\Omega$$



The ideal interferometer Source size

$$R = \int_{Sky} A(\mathbf{s})I(\mathbf{s})\cos(2\pi\nu\mathbf{b}.\mathbf{s}/c) d\Omega$$

= $\cos\left(2\pi\nu\frac{\mathbf{b}.\mathbf{s}_{o}}{c}\right) \int_{Sky} A(\sigma)I(\sigma)\cos(2\pi\nu\mathbf{b}.\sigma/c)d\Omega$
- $\sin\left(2\pi\nu\frac{\mathbf{b}.\mathbf{s}_{o}}{c}\right) \int_{Sky} A(\sigma)I(\sigma)\sin(2\pi\nu\mathbf{b}.\sigma/c)d\Omega$
= $\cos\left(2\pi\nu\frac{\mathbf{b}.\mathbf{s}_{o}}{c}\right) |V|\cos\varphi_{V} - \sin\left(2\pi\nu\frac{\mathbf{b}.\mathbf{s}_{o}}{c}\right) |V|\sin\varphi_{V}$
= $|V|\cos(2\pi\nu\tau_{g} - \varphi_{V})$



The ideal interferometer Summary

• Correlator output:

$$\begin{aligned} r &= \langle E_1 \cos(2\pi\nu t) \ E_2 \cos(2\pi\nu t) \rangle = E_1 \ E_2 \\ r &= E_1 E_2 \cos(2\pi\nu \tau_g) & \longleftarrow \text{ delay} \\ R &= |V| \cos(2\pi\nu \tau_g - \varphi_V) & \longleftarrow \text{ source size} \end{aligned}$$

 \bullet Complex visibility V resembles a Fourier Transform:

$$V = |V|e^{i\varphi_{\rm V}} = \int_{Sky} A(\sigma)I(\sigma)e^{-2i\pi\nu\mathbf{b}.\sigma/c}d\Omega$$

• 3D version of van Cittert–Zernike



The ideal interferometer Bandwidth

• Integrating over a finite bandwidth $\Delta \nu$,

$$R = \frac{1}{\Delta\nu} \int_{\nu_0 - \Delta\nu/2}^{\nu_0 + \Delta\nu/2} |V| \cos(2\pi\nu\tau_g - \varphi_V) \, d\nu$$
$$= |V| \cos(2\pi\nu_0\tau_g - \varphi_V) \, \frac{\sin(\pi\Delta\nu\tau_g)}{\pi\Delta\nu\tau_g}$$

• The fringe visibility is attenuated by a $\sin(x)/x$ envelope (= bandwidth pattern) which falls off rapidly.



The ideal interferometer Bandwidth

$$R = |V| \cos(2\pi\nu_0\tau_g - \varphi_V) \frac{\sin(\pi\Delta\nu\tau_g)}{\pi\Delta\nu\tau_g}$$

- Tracking a source requires **compensation of the geometrical delay** = temporal coherence
- This can be achieved by introducing an **instrumental delay** in the correlator

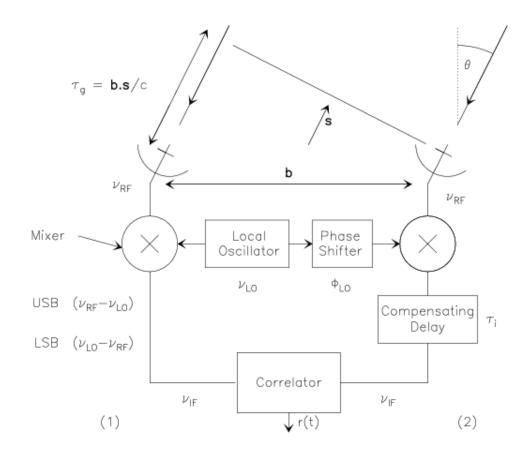


The ideal interferometer Summary

- Correlator output:
 - $\begin{aligned} r &= \langle E_1 \cos(2\pi\nu t) \ E_2 \cos(2\pi\nu t) \rangle = E_1 E_2 \\ r &= E_1 E_2 \cos(2\pi\nu \tau_g) & \longleftarrow \text{ delay} \\ R &= |V| \cos(2\pi\nu \tau_g \varphi_V) & \longleftarrow \text{ source size} \\ R &= |V| \cos(2\pi\nu_0 \tau_g \varphi_V) \frac{\sin(\pi\Delta\nu\tau_g)}{\pi\Delta\nu\tau_g} & \longleftarrow \text{ bandwidth} \end{aligned}$
- We must compensate the delay to obtain $R = |V| \cos(\varphi_V)$



The real interferometer Sketch





• In the receiver **mixer**, the incident electic field is combined with a **local oscillator** signal

$$U(t) = E \cos (2\pi\nu t + \varphi)$$

$$U_{\rm LO}(t) = E_{\rm LO} \cos (2\pi\nu_{\rm LO}t + \varphi_{\rm LO})$$

$$\nu_{\rm LO} \simeq \nu$$

• The mixer is a **non-linear** element:

 $I(t) = a_0 + a_1(U + U_{\rm LO}) + a_2(U + U_{\rm LO})^2 + a_3(...)^3 + ...$



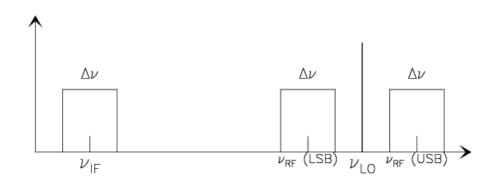
- There are terms at various frequencies and harmonics
- A **filter** selects the frequencies such that;

 $\nu_{\rm IF} - \Delta \nu / 2 \le |\nu - \nu_{\rm LO}| \le \nu_{\rm IF} + \Delta \nu / 2$

- $\nu_{\rm IF}$ is the **intermediate frequency**
- $\nu_{\rm IF}$ such that amplifiers and transport elements available
- \bullet PdBI: $\nu_{\rm IF}=4\text{--}8$ GHz, ALMA: $\nu_{\rm IF}$ =4-12 GHz

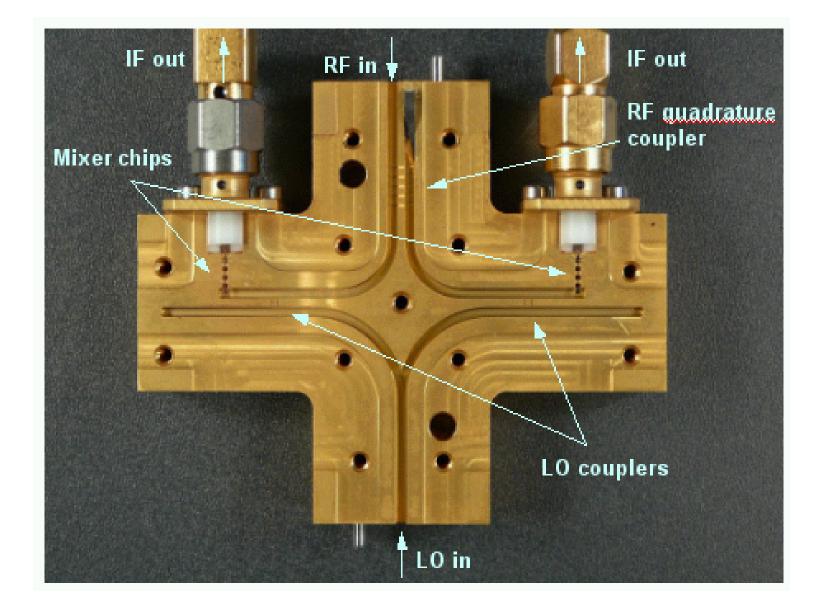


- The receiver output is
 - $I(t) \propto E E_{\rm LO} \cos \left(\pm \left(2\pi (\nu \nu_{\rm LO}) t + \varphi \varphi_{\rm LO} \right) \right)$



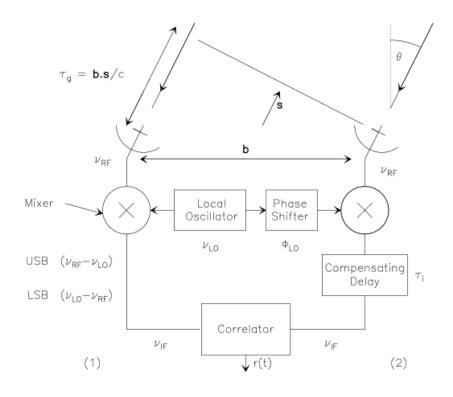


- **DSB** receivers accept both LSB and USB frequencies, i.e. their output is the sum of LSB and USB
- **SSB** receivers accept only LSB or USB (response very strongly frequency dependant)
- **2SB** receivers are 2 DSB receivers combined such that the two bands are independently output (and processed)





The real interferometer Delay tracking



• A compensating delay is introduced in one of the branch of the interferometer, **on the IF signal**

• Equivalent to the delay lines in IR in-terferometers



The real interferometer Delay tracking

• Phases of the two signals (USB):

$$\varphi_1 = 2\pi\nu\tau_g \quad \varphi_1 = 2\pi\nu\tau_g = 2\pi(\nu_{\rm LO} + \nu_{\rm IF})\tau_g$$

$$\varphi_2 = 0 \qquad \varphi_2 = 2\pi\nu_{\rm IF}\tau_i$$

• Correlator output:

$$R = |V| \cos(2\pi\nu\tau_g - \varphi_V)$$
$$R = |V| \cos(\varphi_1 - \varphi_2 - \varphi_V)$$
$$R = |V| \cos(2\pi\nu_{\rm LO}\tau_g - \varphi_V)$$



The real interferometer Fringe Stopping

- \bullet Delay tracking not enough because applied on the IF
- Solution: in addition to delay tracking, **rotate the phase of the local oscillator** such that at any time:

$$\varphi_{\rm LO}(t) = 2\pi\nu_{\rm LO}\tau_g(t)$$

- τ_g is computed for a reference position = phase center
- Phase center = pointing center in practice, though not mandatory



The real interferometer Fringe stopping

• Phases of the two signals (USB):

$$\varphi_{1} = 2\pi\nu\tau_{g} = 2\pi(\nu_{\rm LO} + \nu_{\rm IF})\tau_{g}$$
$$\varphi_{2} = 2\pi\nu_{\rm IF}\tau_{i} + \varphi_{\rm LO}$$
$$\varphi_{\rm LO} = 2\pi\nu_{\rm LO}\tau_{g}$$

• Correlator output:

$$R = |V| \cos(\varphi_1 - \varphi_2 - \varphi_V)$$
$$R = |V| \cos(\varphi_V)$$



The real interferometer Complex correlator

• After fringe stopping:

$$r_r = |V|\cos(-\varphi_{\rm V})$$

- No time/delay dependance any more \longrightarrow cannot measure |V| and $\varphi_{\rm V}$ separately.
- A second correlator is necessary, with one signal phase shifted by $\pi/2$:

 $r_i = |V|\sin(-\varphi_{\rm V})$

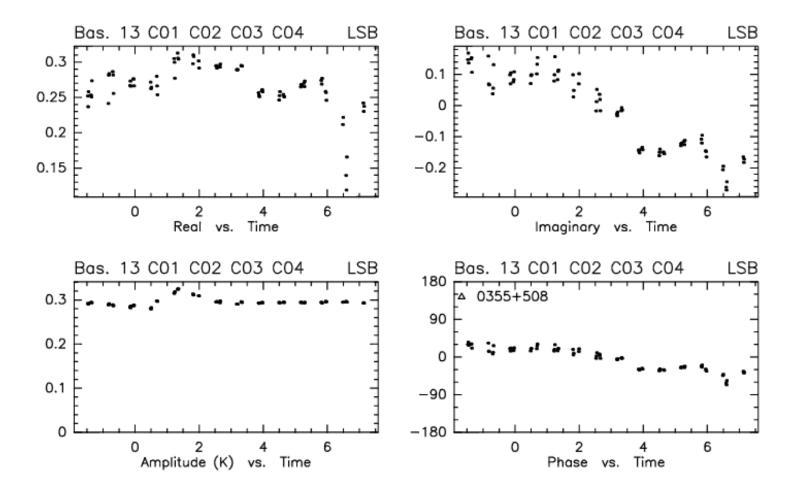
• The complex correlator measures directly the visibility



The real interferometer Complex correlator

- The correlator measures the real and imaginary parts of the visibility. **Amplitude and phases are computed off-line.**
- Amplitude and phases have more physical sense
 - -Visibility amplitude = **correlated flux**
 - The atmosphere adds a **phase** to the incoming signals \longrightarrow measured phase = visibility + $\varphi_1 - \varphi_2$

RF:	Uncal.	CLIC - 06-0CT-2008 11:19:29 - boissier@pctcp04 W08E03W05N02N07 6Dq-N11	Scan Avg.
Am:	Abs.	R9 HCN(1-0) 88.782GHz B1 Q3(320,320,320,20)V Q3(320,320,320,20)H	Narrow Input 1
Ph:	Abs.	(182 2942 P CORR)-(981 3562 P CORR) 26-0CT-2007 22:31-07:09	





The real interferometer Spectroscopy

- Remember the Wiener-Kichnine theorem?
- Calculate the correlation function for several delay $\delta \tau \longrightarrow$ measurement of the **temporal correlation** \longrightarrow FT to get the spectra:

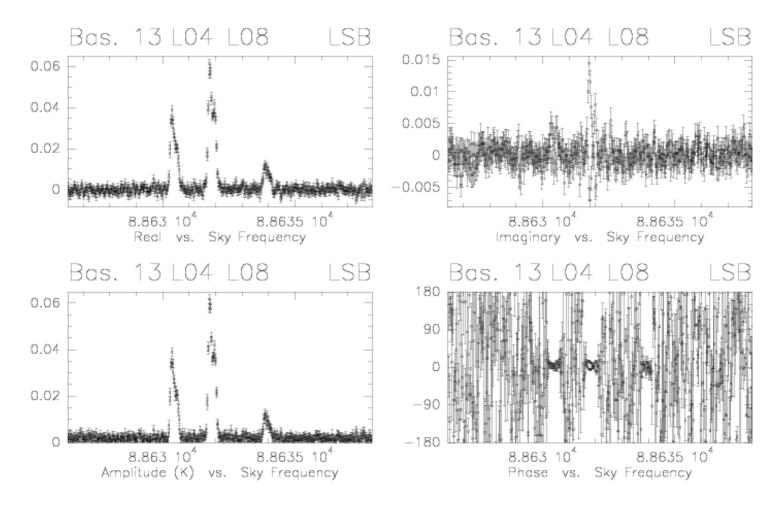
$$V_{\nu}(u,v,\nu) = \int V(u,v,\tau) e^{-2i\pi\tau\nu} d\nu$$

- Nothing to do with geometrical delay compensation $\delta \tau \sim 1/\delta \nu$ here
- Mixed up implementation in correlator software

 RF:
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 CLIC - 06-0CT-2008 09:54:09 - boissier@pctcp04
 W08E03W05N02N07 6Dq-N11
 Scan Avg.

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 BOTH palarizations

 Ph:
 Abs.
 (146 2909 0 CORR)-(972 3556 0 CORR) 26-0CT-2007 22:07-07:05
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Goal Interferometry

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Aperture synthesis Complex visibility

• Complex visibility:

$$V = |V|e^{i\varphi_{\rm V}} = \int_{Sky} A(\sigma)I(\sigma)e^{-2i\pi\nu\mathbf{b}.\sigma/c}d\Omega$$

- Going from 3-D to 2-D? ...some algrebra...
- OK providing that: (max. field of view)² × max. baseline $\ll 1$ $\implies \frac{(\max. \text{ field of view})^2}{\text{resolution}} \ll 1$



Aperture synthesis Complex visibility

$$V(u,v) = \int_{Sky} A(\ell,m) I(\ell,m) e^{-2i\pi\nu(u\ell+vm)} d\Omega$$

- *uv* plane is perpendicular to the source direction, fixed
 wrt source → back to Young's hole
- Price: limit on the field of view
- Approximation **ok in (sub)mm domain**, problem at wavelengths > cm



Aperture synthesis (Field of view)

- Field of view is limited by
 - the **antenna primary beam**: the interferometer measures $A \times I$
 - the 2D visibility approximation
 - the frequency averaging (bandwidth)
 - the time averaging (integration)
 - \hookrightarrow averaging in the uv plane; possible only if limited field of view



Aperture synthesis (Field of view)

• Values for Plateau de Bure

0.5 GHz 1 Min $\theta_{\rm s}$ Primary 2-1) \mathcal{V} (GHz) Field Bandwidth Averaging Beam 5″ 5' 80″ 60" 80 2' 2" 80 3.5' 30" 45^{'''} 60″ 2" 45^{'''} $230 \quad 3.5' \quad 1.5'$ 24″ 22" 1.7' 12" 24″ 0.5''230

- Problem with 2D field: software; with bandwith: split the data for imaging; with time averaging: dump faster.
- Primary beam is the main limit on the FOV

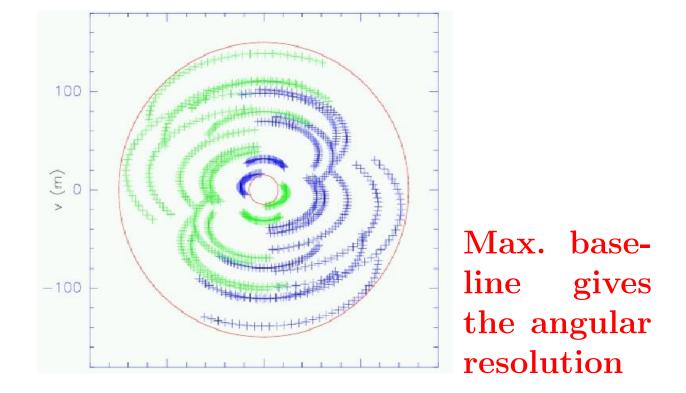


Aperture synthesis uv plane

- *uv* plane is perpendicular to the source direction, fixed
 wrt source → back to Young's hole
- (u, v) is the 2-antennas **vector** baseline projected on the plane perpendicular to the source
- (u, v) are spatial frequencies
- ... Earth rotation ... (spherical trigonometry) ...
- (u, v) describe an **ellipse** in the uv plane (for $\delta = 0 \deg$, a line)



Aperture synthesis uv plane coverage





Aperture synthesis Transfer function

Single-dish observations

Aperture function \Rightarrow Voltage pattern $\downarrow |\cdot|^2$ \Rightarrow Power pattern $B(\ell, m)$ = **Primary beam**



Aperture synthesis Transfer function

Single-dish observations

Transfer function describes how spatial frequencies are transmitted by the telescope



Aperture synthesis Transfer function

Interferometers

Aperture function \Rightarrow Voltage pattern $\star \downarrow$ $\downarrow |\cdot|^2$ Transfer function $\mathbf{T}(\mathbf{u}, \mathbf{v})$ \Rightarrow Power pattern $B(\ell, m)$ = Primary beam

Aperture synthesis = sample directly the transfer function



Sensitivity Radiometric formula

- Measurement of visibilities is limited by noise emitted by atmosphere, antenna, ground, receivers.
- The rms noise for the baseline ij is given by:

$$\delta S_{ij} = \frac{2k}{A\eta_{\rm A}\eta_{\rm Q}\eta_{\rm P}} \cdot \frac{\sqrt{T_{\rm SYSi}T_{\rm SYSj}}}{\sqrt{2B\,T}}$$

-B bandwidth

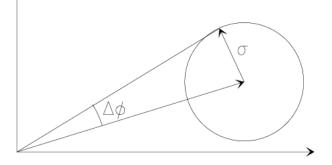
-T integration time

- $-\,A$ antenna physical aperture
- $\, \eta_{\scriptscriptstyle \rm A}$ antenna aperture efficiency
- $\, \eta_{\scriptscriptstyle \mathrm{Q}}$ efficiency for the correlator
- $-T_{\text{sys}i}$ system noise temperature (single dish)
- $-\eta_{\rm P}$ phase decorrelation factor (LO jitter) le dish)



Sensitivity Radiometric formula

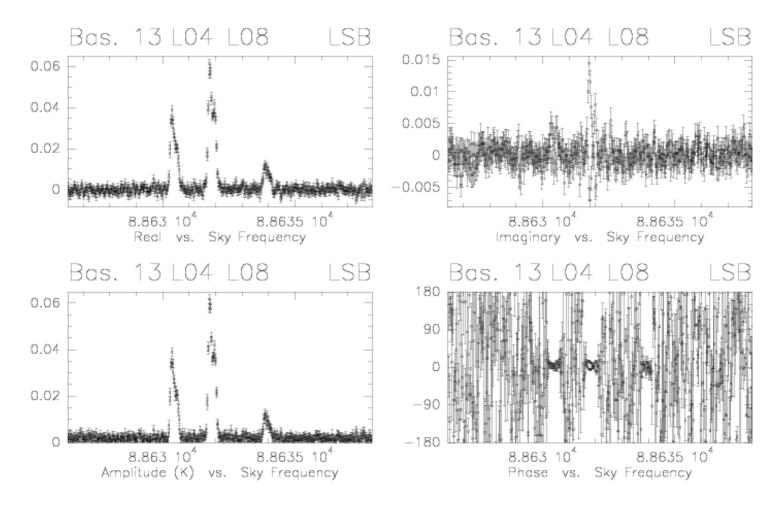
- This is the noise on the **real** and on the **imaginary** parts of the visibilities (measured independently)
- \bullet This is also the noise on the **amplitude** S
- \bullet Noise on the phase more complex, of the order of σ/S



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 (146 2909 0 CORR)-(972 3556 0 CORR) 26-0CT-2007 22:07-07:05
 BOTH palarizations





Sensitivity Radiometric formula

• For N identical antenna/receivers, i.e. N(N-1)/2 baselines, the **point-source** sensitivity is:

$$\delta S = \frac{2k}{A\eta_{\rm A}\eta_{\rm Q}\eta_{\rm P}} \cdot \frac{T_{\rm SYS}}{\sqrt{N(N-1)\,B\,T}}$$

- Scales as $\sim 1/N$
- Sensitivity to extended sources depends on angular resolution



Sensitivity Phase decorrelation

• Short term phase errors in the local oscillators (jitter) will cause a **decorrelation** of the signal and reduce the visibility amplitude by a factor

$$\eta_{\rm P}(12) = e^{-(\sigma_1^2 + \sigma_2^2)/2} = \eta_1 \eta_2$$

• Requirements:

$$\begin{array}{c|ccccc} \eta_1 & 0.99 & 0.98 & 0.95 & 0.90 \\ \hline \sigma_1 \ (\text{degrees}) & 8.1 & 11.5 & 18.3 & 26.4 \end{array}$$



Sensitivity Phase decorrelation

- $\eta_{\rm P} = 0.9 \longrightarrow \eta_1 = 0.95 \longrightarrow \sigma_1 = 18 \deg$
- PdBI: LO derived from a reference at 1.8 GHz
- Phase stability required = $\sigma_1 (1.8 \text{ GHz}/230 \text{ GHz}) \sim 0.15^{\circ}$
- Very stable oscillators are required
- Phase decorrelation due to the atmosphere more severe problem



Summary Other instrumental issues

- Antenna position measurements, to get the delay, u, v
- Phase lock systems to control $\varphi_{\rm LO}$
- Real-time monitoring and correction of the phase offset in the cables or fibers
- Antenna deformations, e.g. thermal expansion (delay)
- Accurate focus measurements (delay)
- Atmospheric phase monitoring





Summary It works!

