

Operation of resonators in the nonlinear regime

Loren Swenson

Resonator Workshop

July 28, 2011



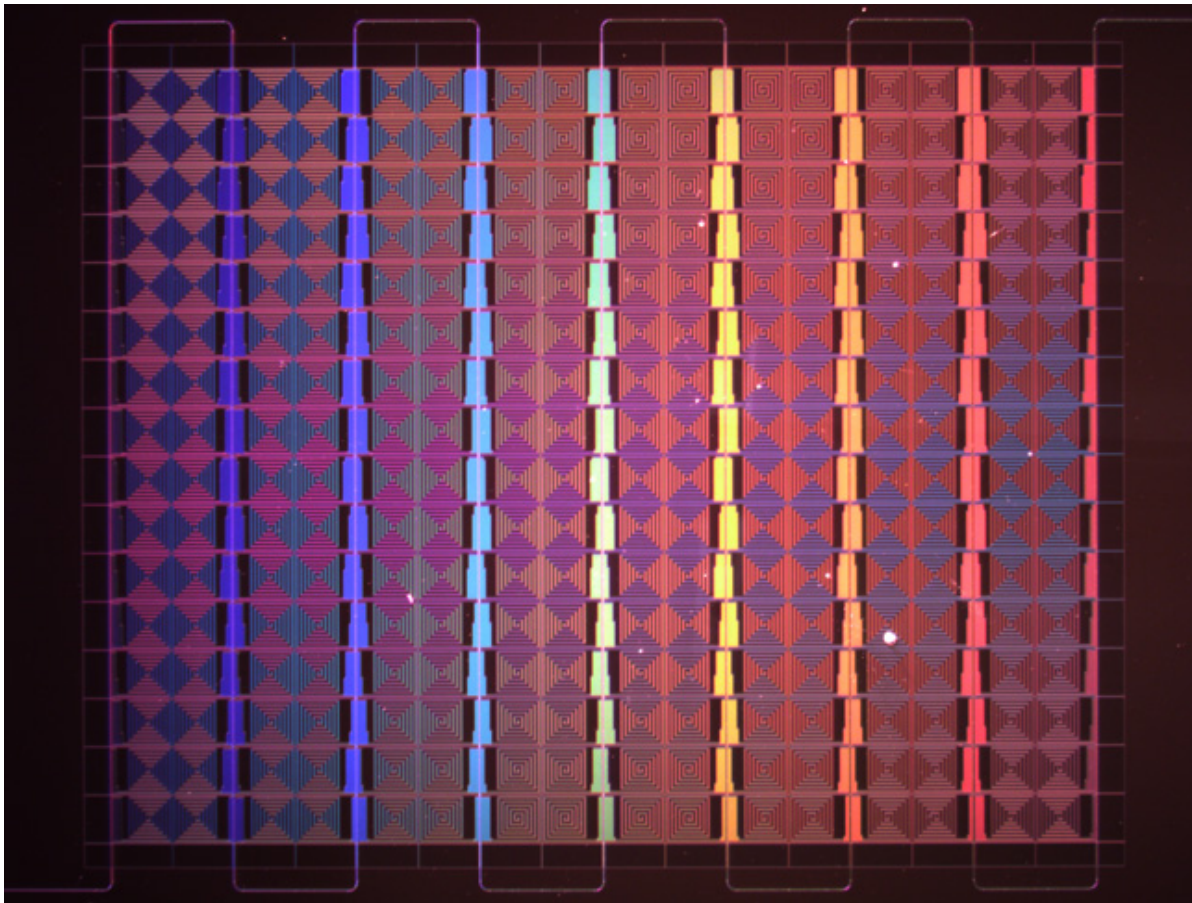
Caltech



Motivation

TiN is a very promising material:

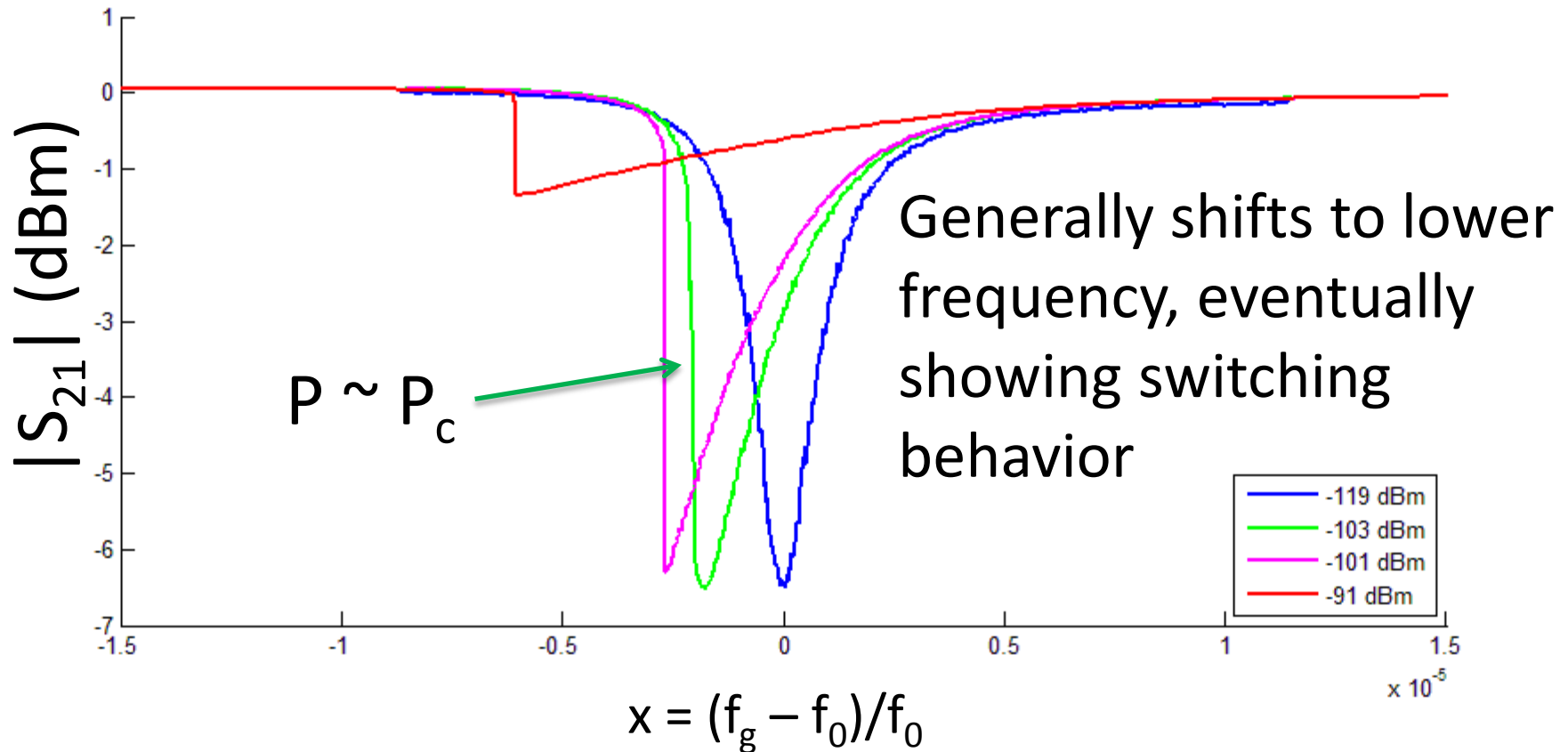
- High normal state resistivity (70-200 $\mu\Omega\text{-cm}$)
- Internal Qs exceeding 10^7
- Tunable T_c
- High kinetic inductance fraction



- 16x16 LEKID array. 100 nm TiN on Si with $T_c \sim 2$ K.
- Optimized for absorbing FIR radiation.
- Band-defining filters + Blackbody used to illuminate with ~ 5 pW radiation
- $Q_c \sim Q_i = 800,000$; $Q_r \sim 400,000$

However: Onset of nonlinearity observed at low powers

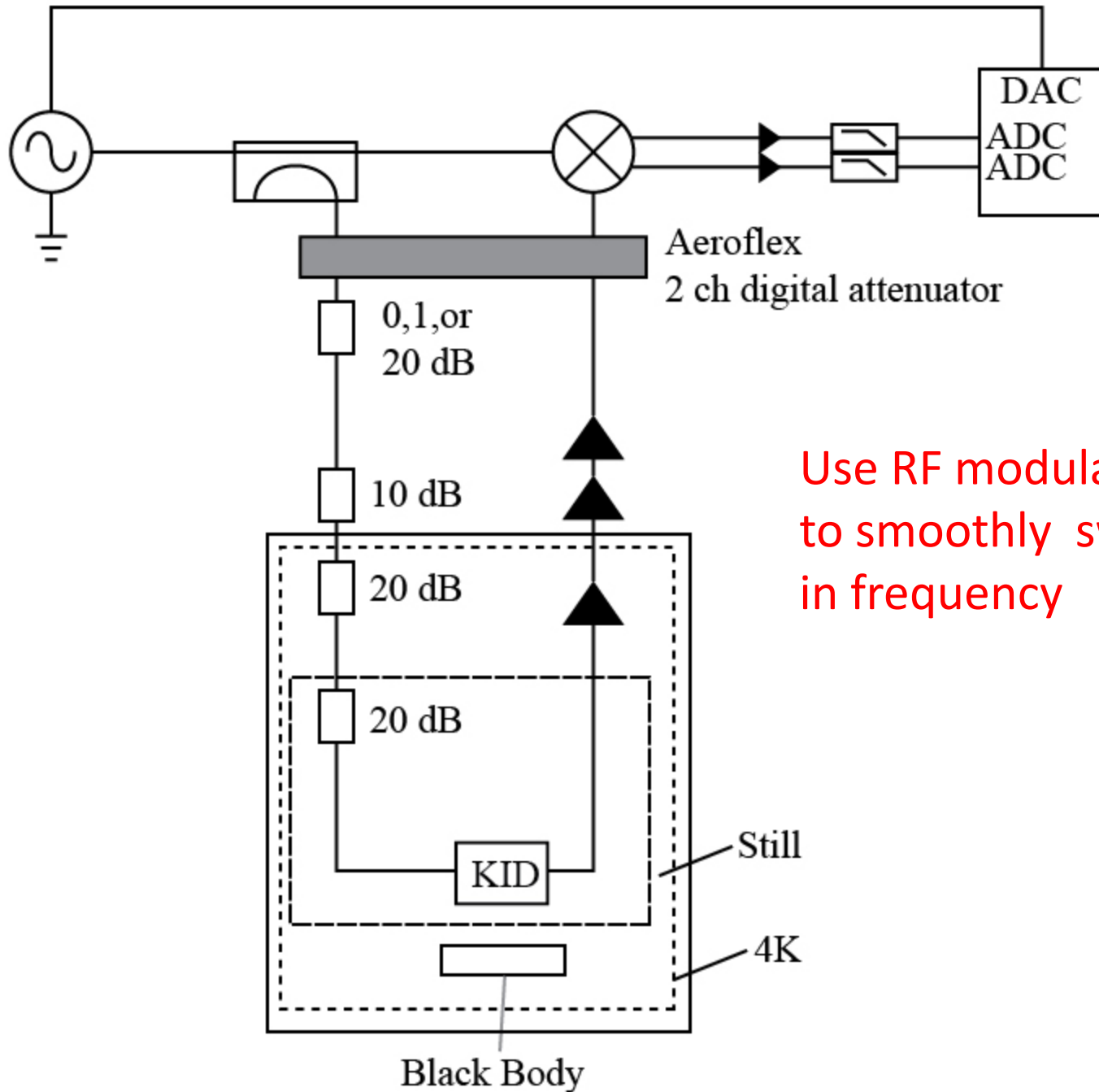
Typical upward sweeping VNA scans



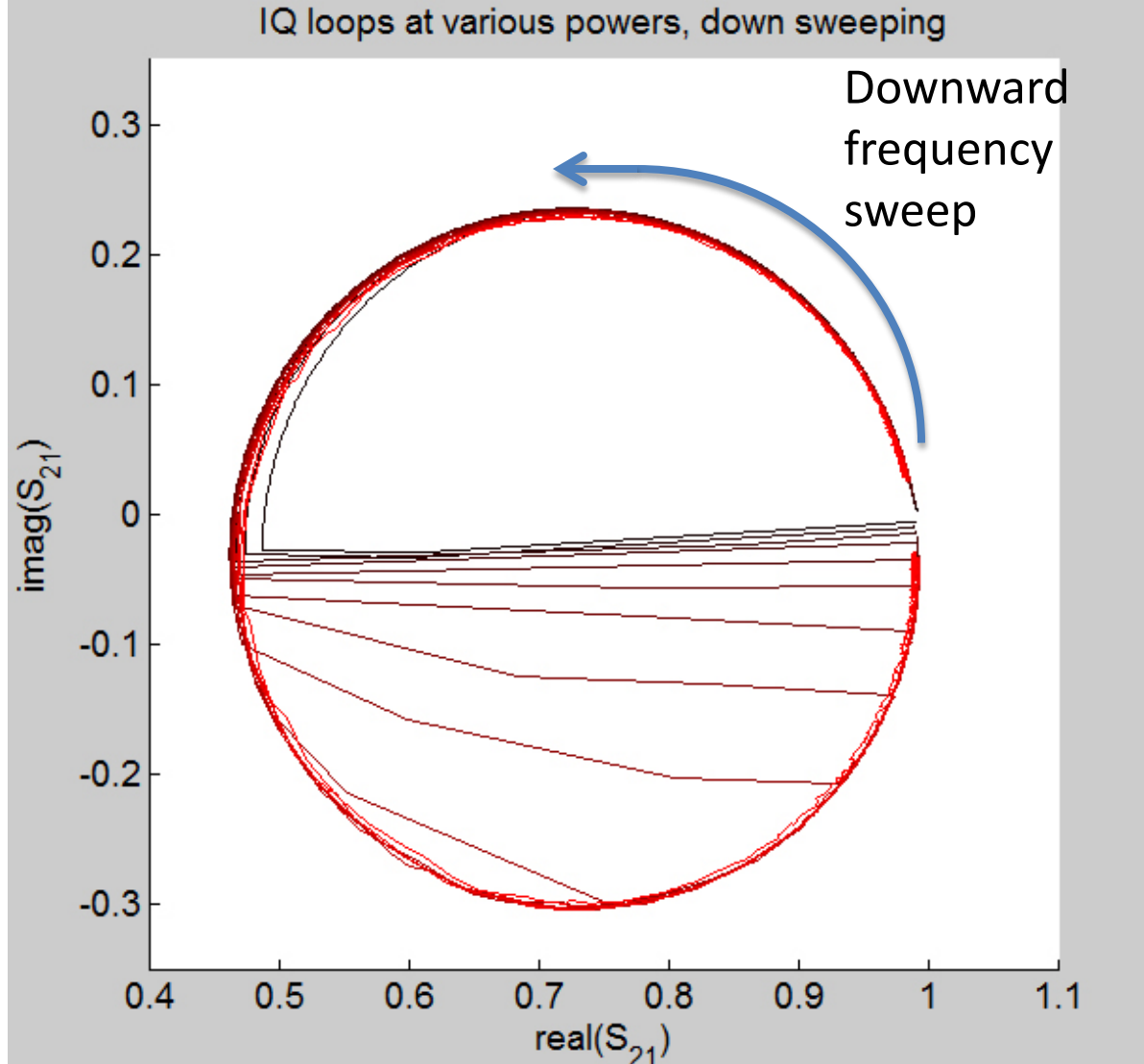
Higher power operation is desirable.

- 1) Suppress amplifier noise in dissipation and frequency direction.
- 2) Suppress TLS noise in the frequency direction.

Question: Can we understand the nonlinear behavior? Does this allow us to operate at higher powers?



Use RF modulated source to smoothly sweep downward in frequency



- Diameter of the IQ loop doesn't change significantly over broad power range (35 dB).
- Both well below and above onset of bifurcation.

- Q_i observed to be fairly constant over large power range, above and below bifurcation.
- Kinetic inductance known to exhibit nonlinear behavior at large currents
- Due to symmetry considerations, leading term quadratic:

$$L_{\text{kin}}(\mathbf{I}) = L_{\text{kin}}(0)[1 + \mathbf{I}^2/\mathbf{I}_*^2]$$

Nonlinear scale factors

$$E_* \propto L\mathbf{I}_*^2$$

$$P_* = w_r E_* / Q_r$$

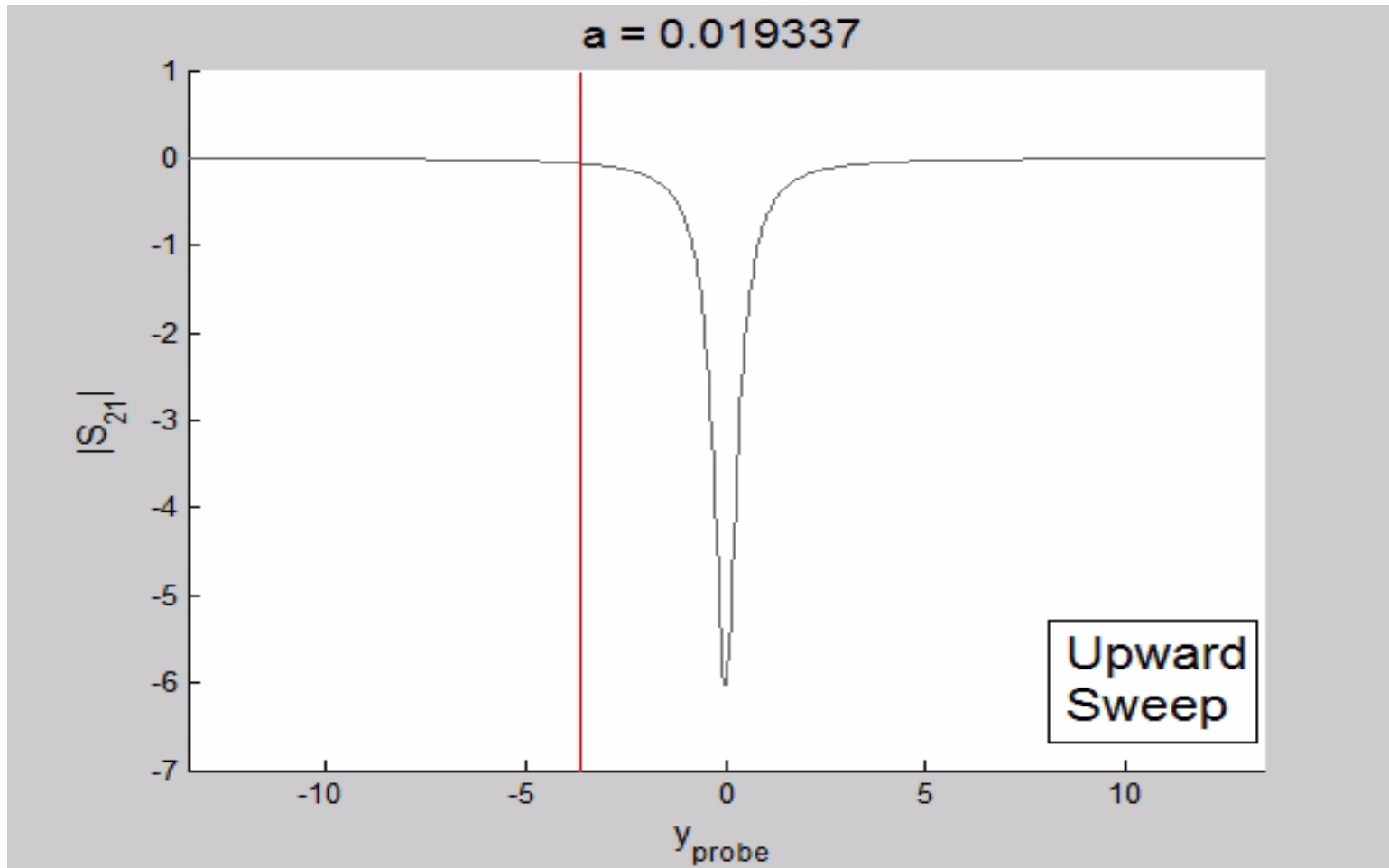
$$a = (\chi_c Q_r / 2)^* (P / P_*)$$

$$a_c \approx .77$$

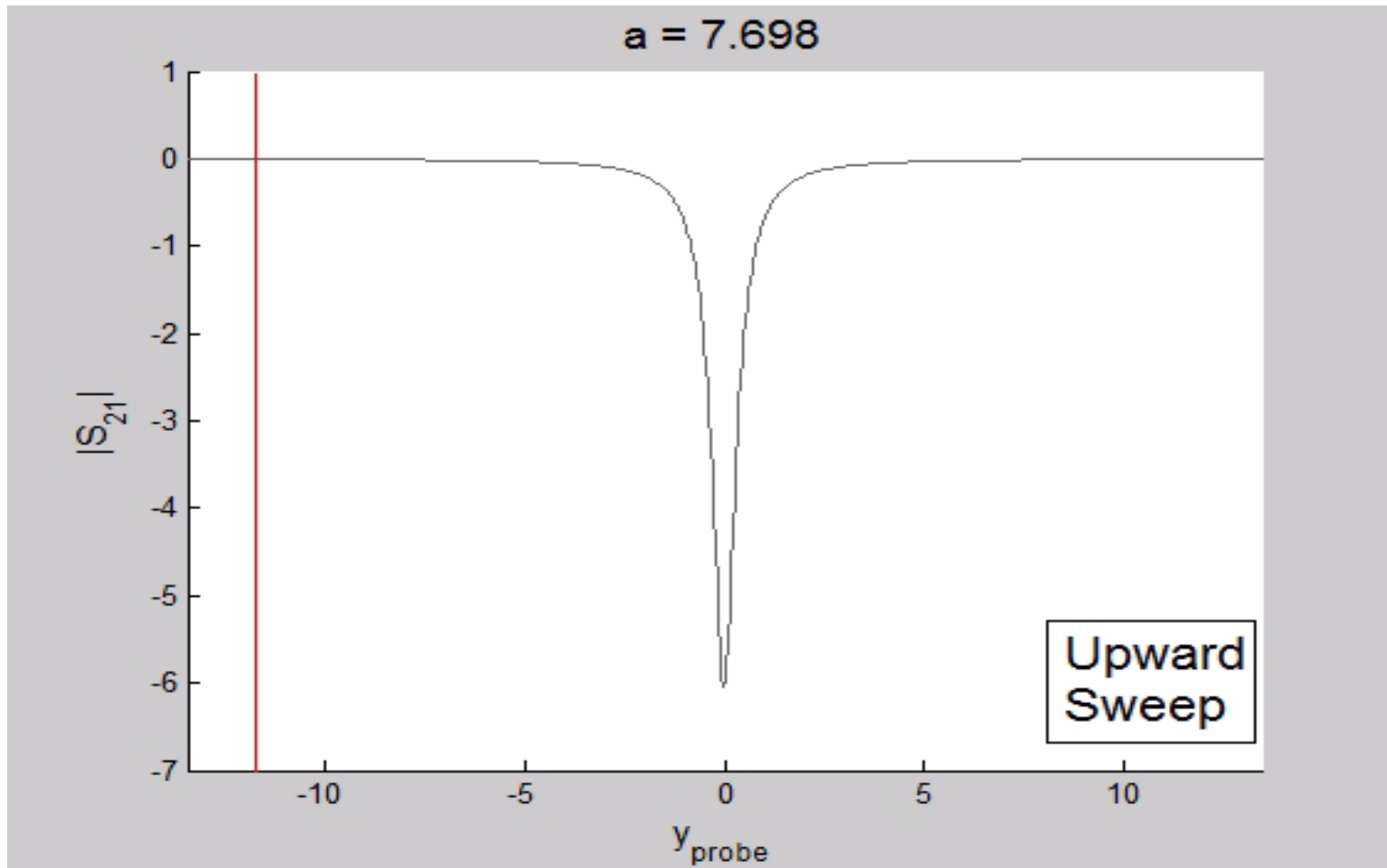
Fractional frequency shift measured in linewidths

$$y_g = (f_g - f_{r,0}) / \Delta f$$

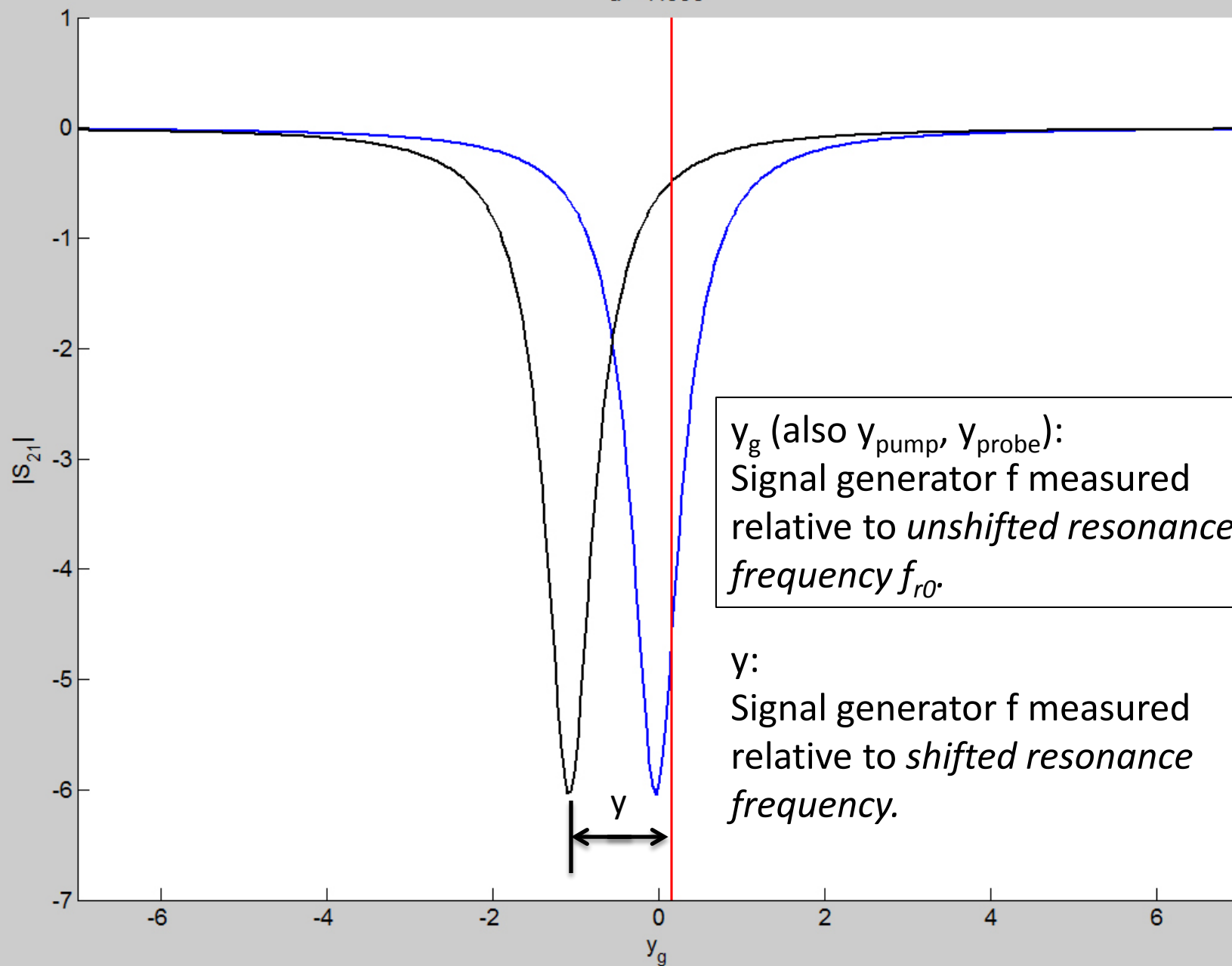
VNA measurement:
Probe power well **below** the onset of nonlinearity



VNA measurement:
Probe power well **above** the onset of nonlinearity



$a = 7.698$



y_g (also y_{pump} , y_{probe}):
Signal generator f measured
relative to *unshifted resonance*
frequency f_{r0} .

y :
Signal generator f measured
relative to *shifted resonance*
frequency.

Using conservation of energy, possible to derive the expression:

$$y_g = \mathbf{y} - a/(1+4\mathbf{y}^2)$$

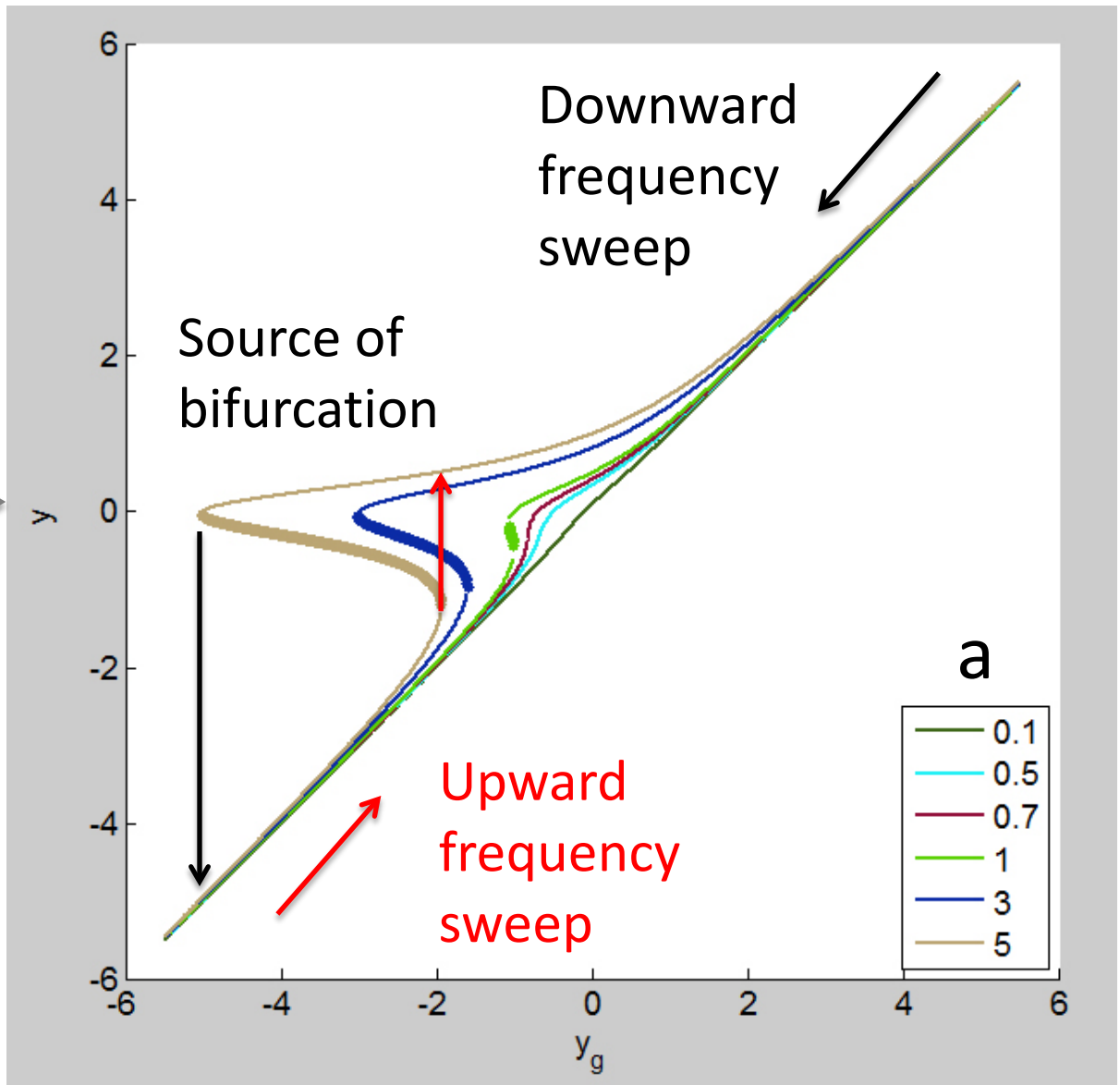
Solve for $y(y_g)$



Why this matters to everyone (even people operating at $P < P_c$):

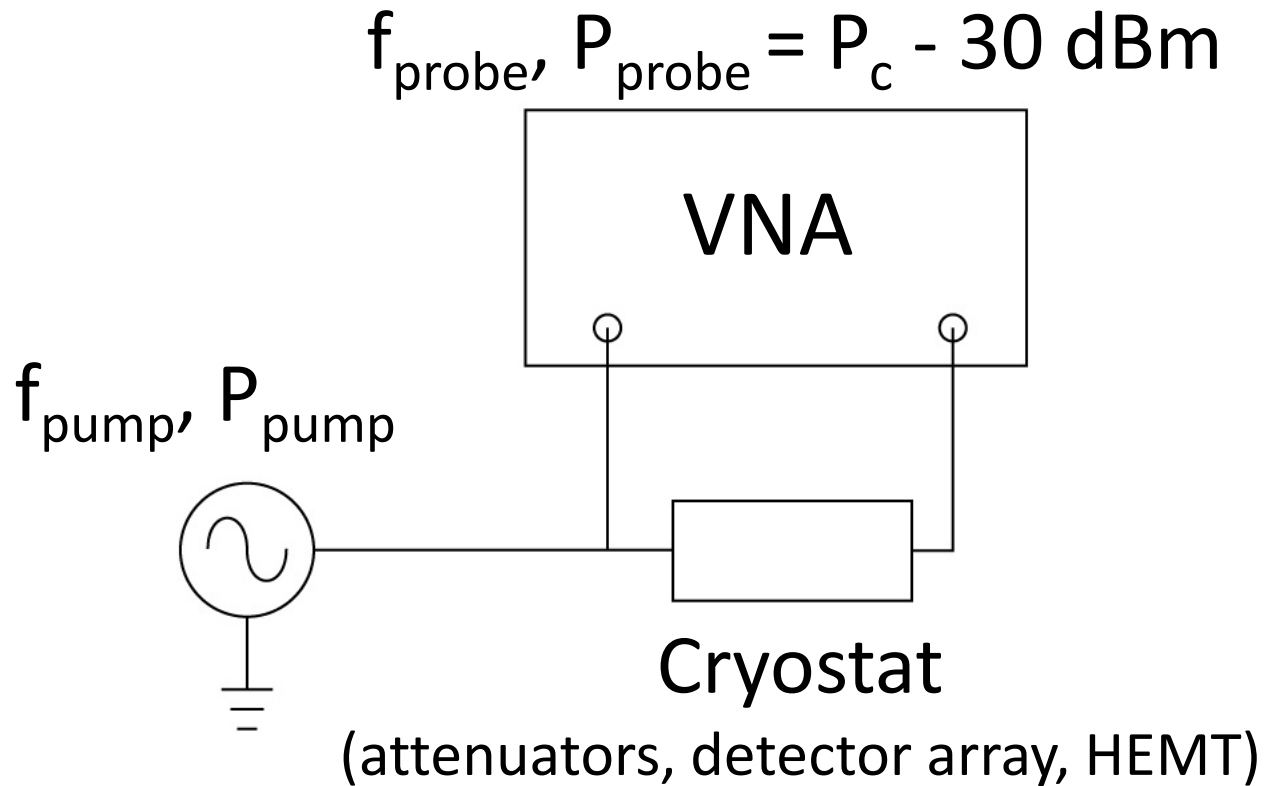
$$S_{21} = 1 - Q_r/Q_c * [1/(1+2j\mathbf{y})]$$

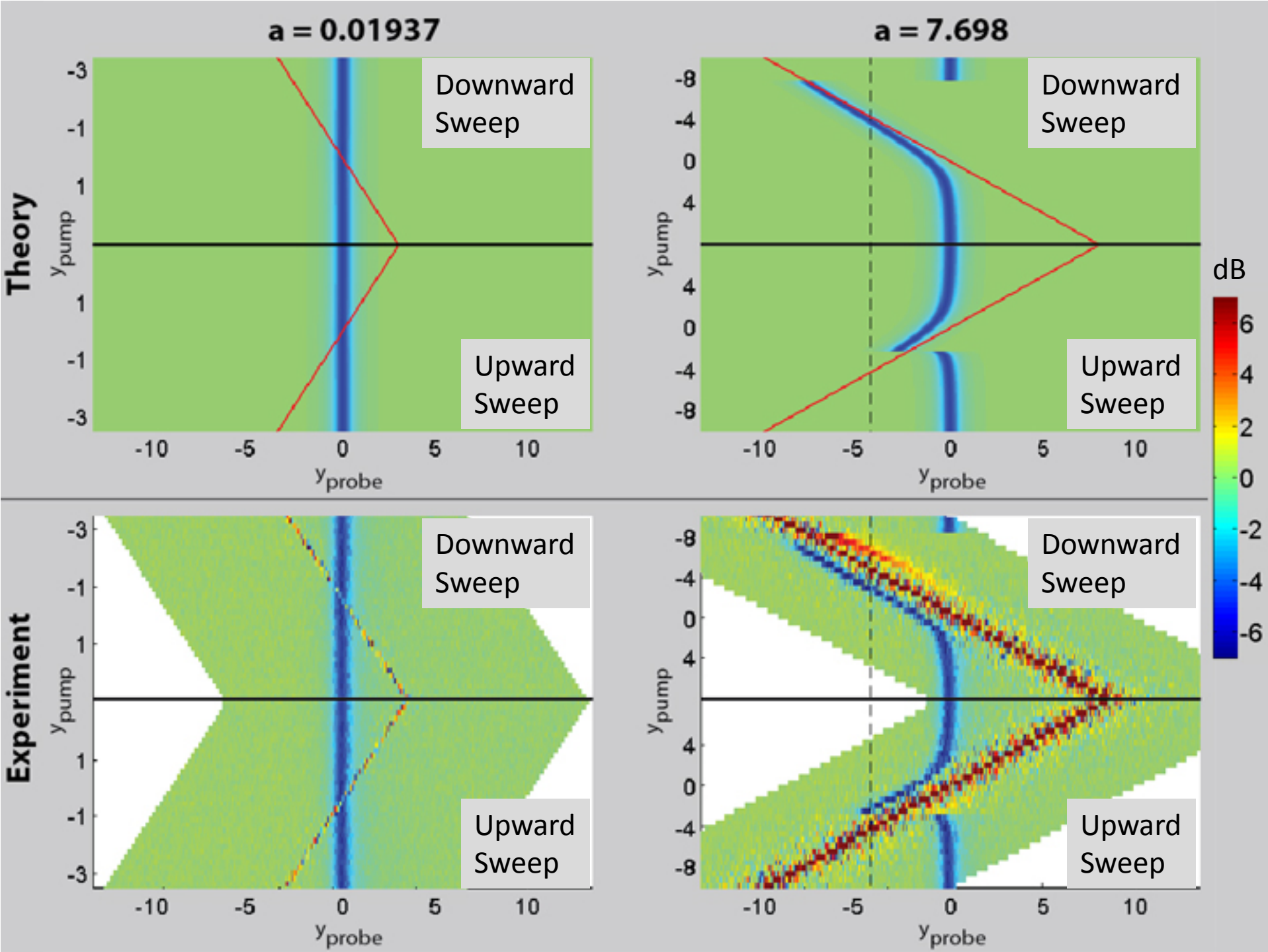
(We'll come back to this in a few slides)



Is this really what is going on?

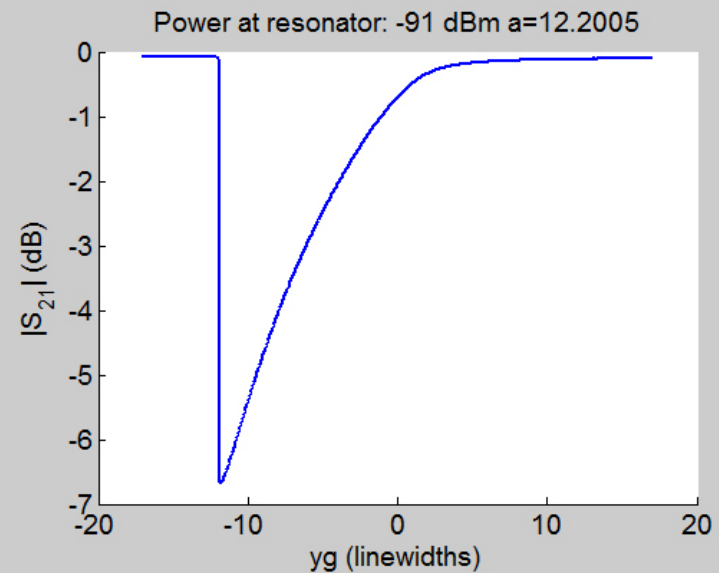
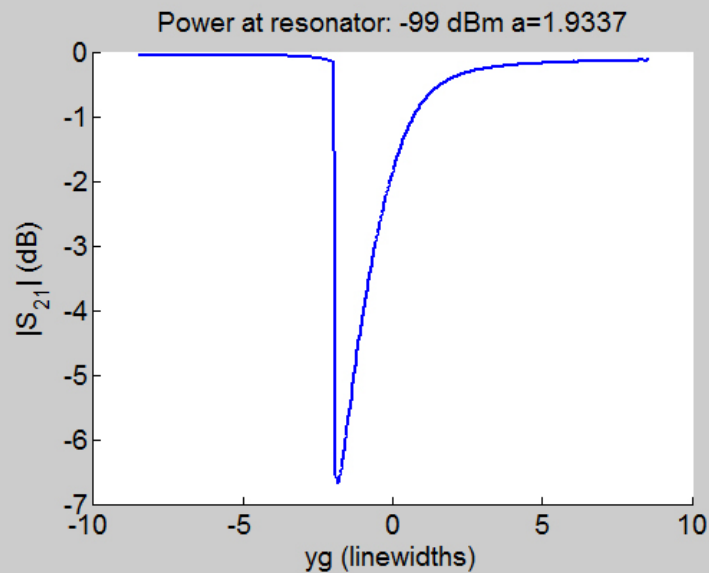
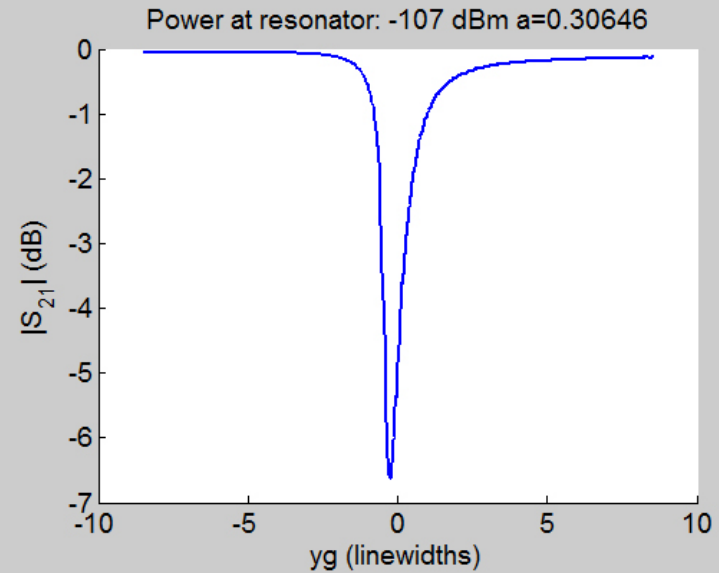
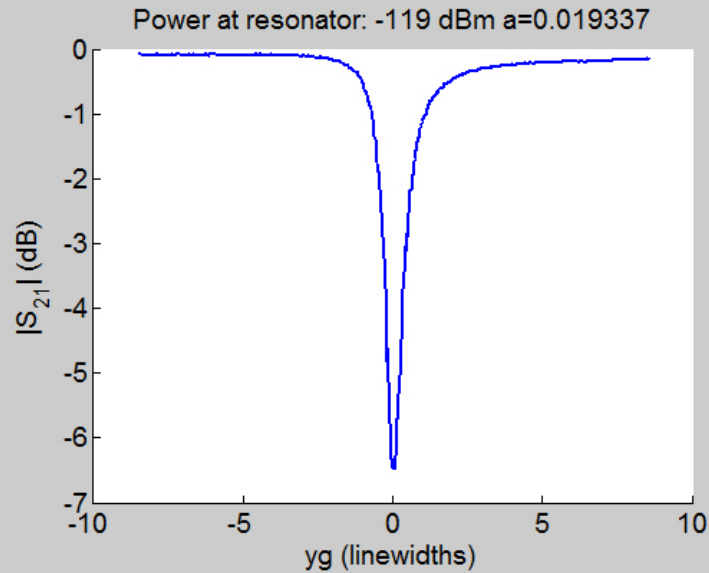
ie Can we measure this experimentally?





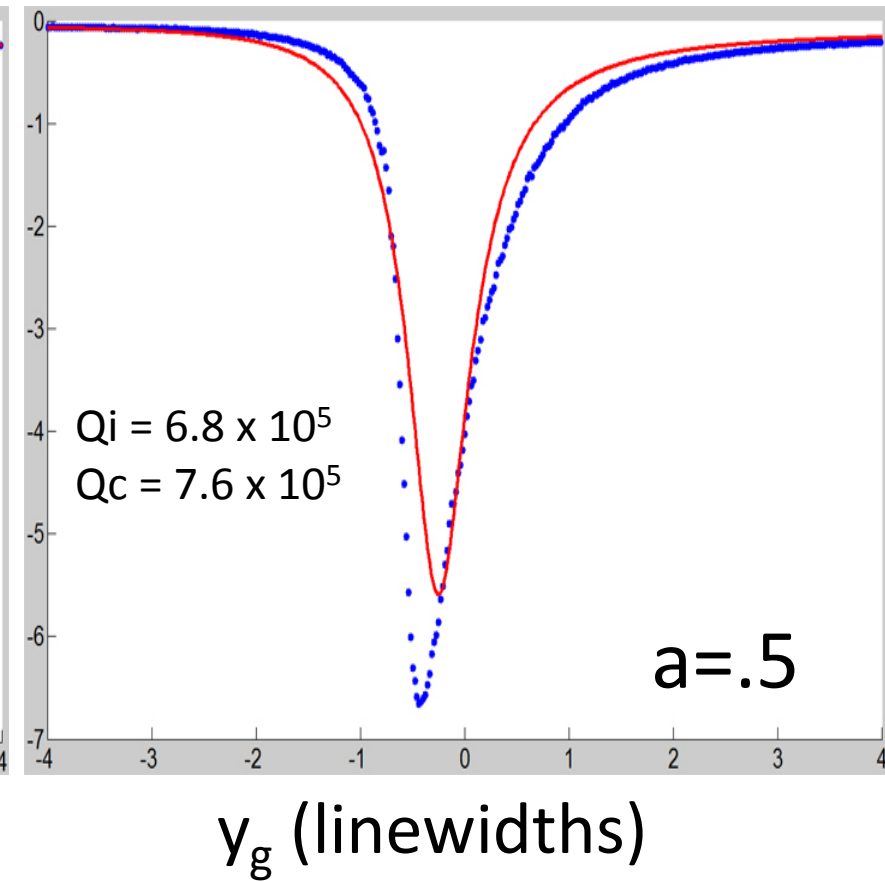
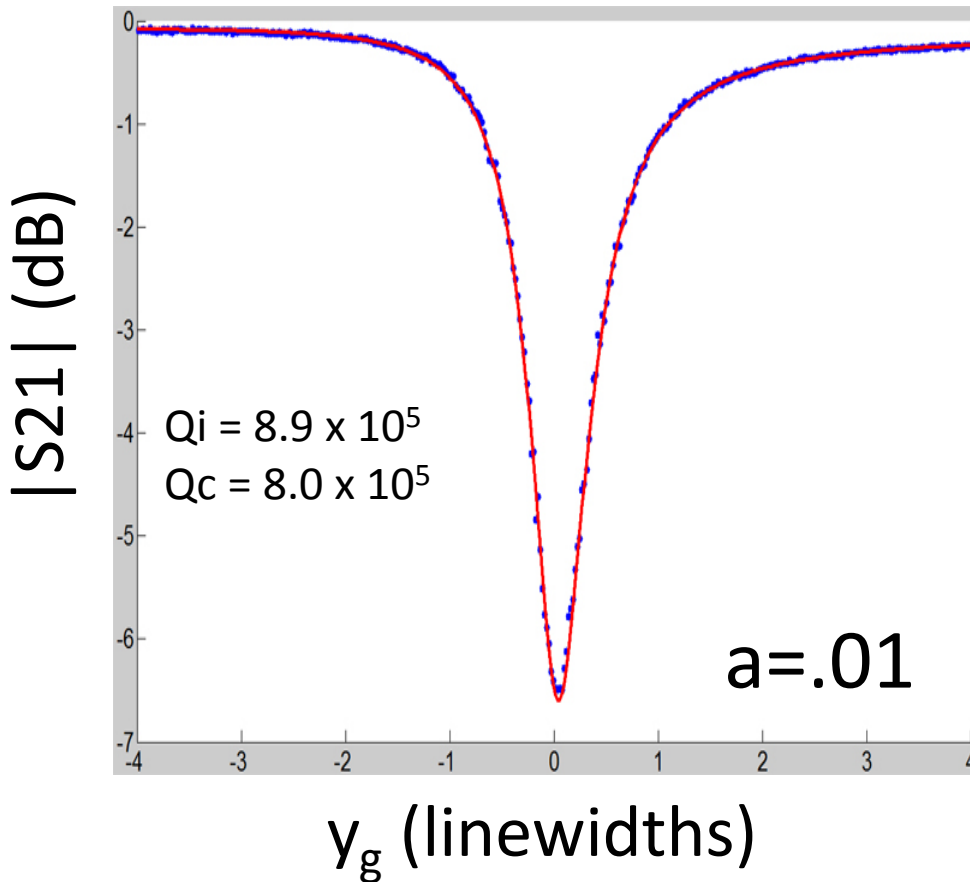
Fitting nonlinear resonances

$$S_{21} = 1 - Q_r/Q_c * [1/(1+2jQ_r x)]$$



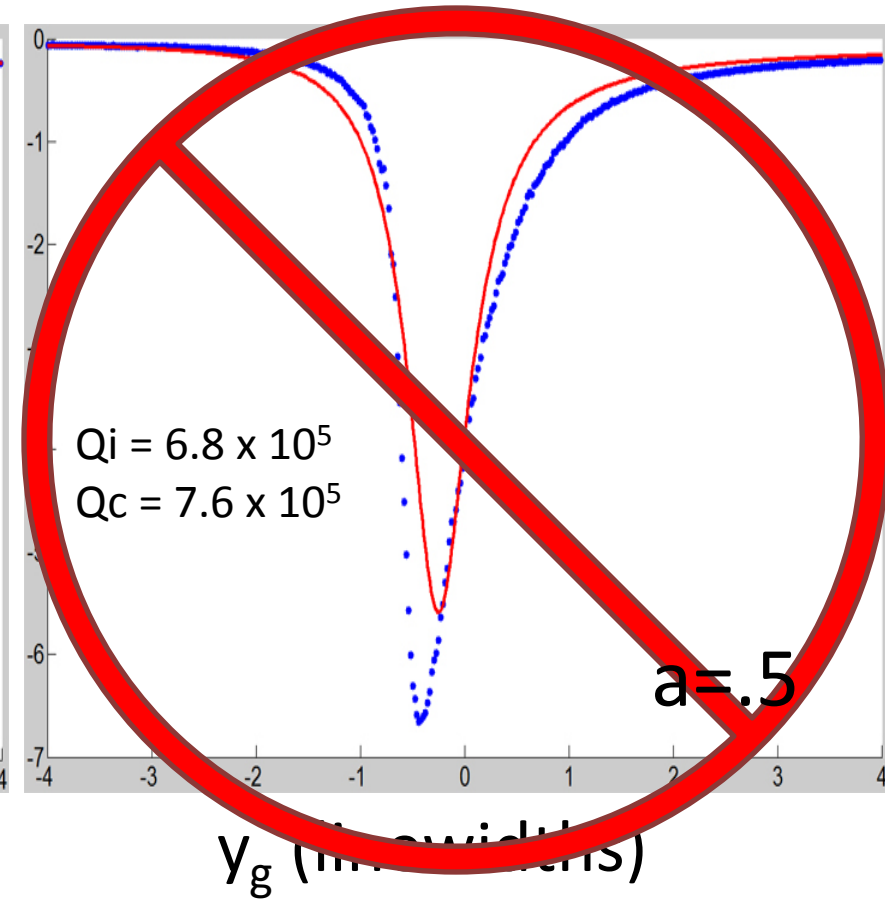
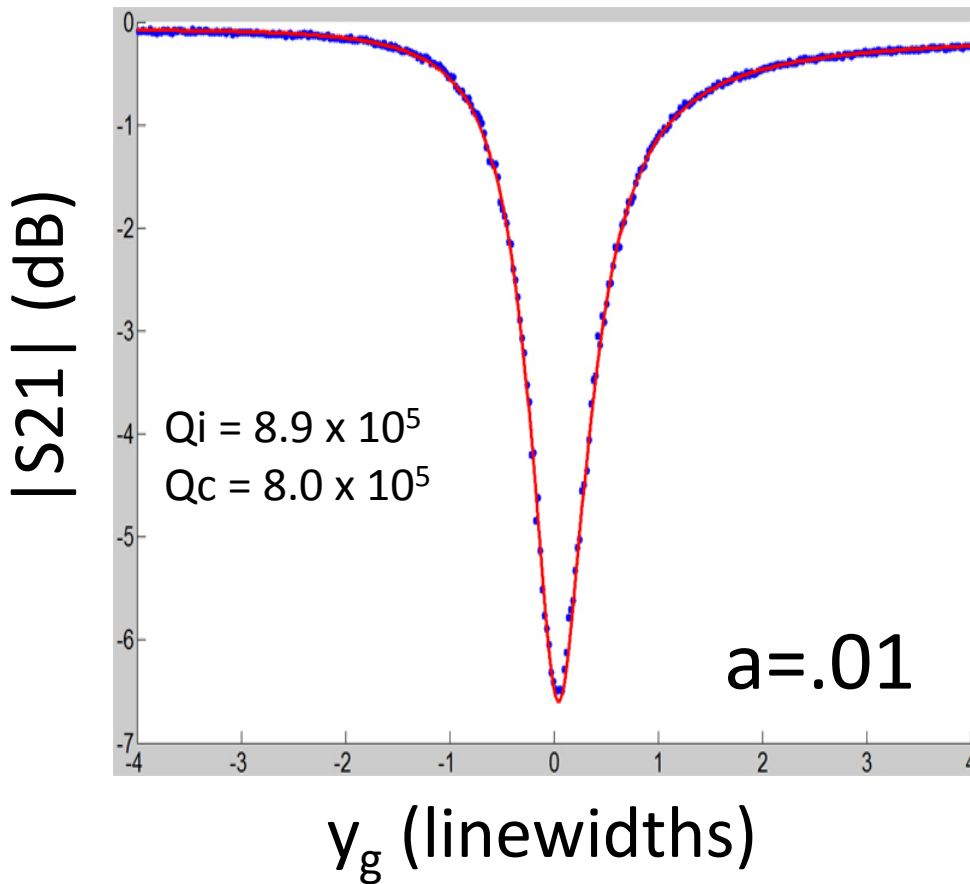
Fitting single-tone curves

$$t_{21}(f) = ae^{-2\pi jf\tau} \left[1 - \frac{Q_r/Q_c e^{j\phi_0}}{1 + 2jQ\left(\frac{f-f_r}{f_r}\right)} \right]$$



Fitting single-tone curves

$$t_{21}(f) = ae^{-2\pi jf\tau} \left[1 - \frac{Q_r/Q_c e^{j\phi_0}}{1 + 2jQ\left(\frac{f-f_r}{f_r}\right)} \right]$$



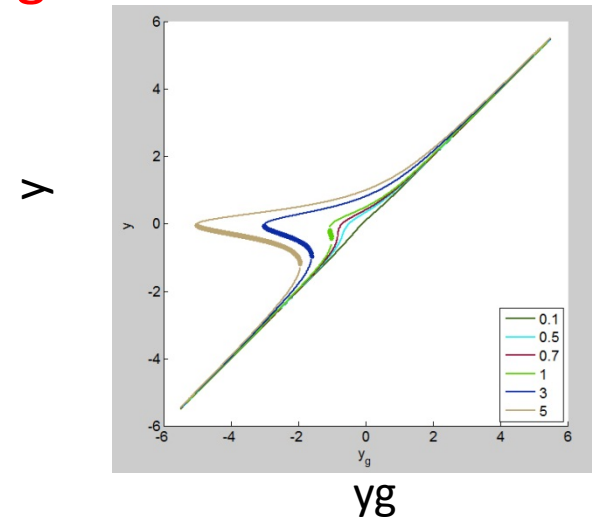
Fitting single-tone curves

$$t_{21}(f) = ae^{-2\pi jf\tau} \left[1 - \frac{Q_r/Q_c e^{j\phi_0}}{1 + 2j\gamma} \right]$$

$$\gamma_g = \gamma - a/(1+4\gamma^2)$$

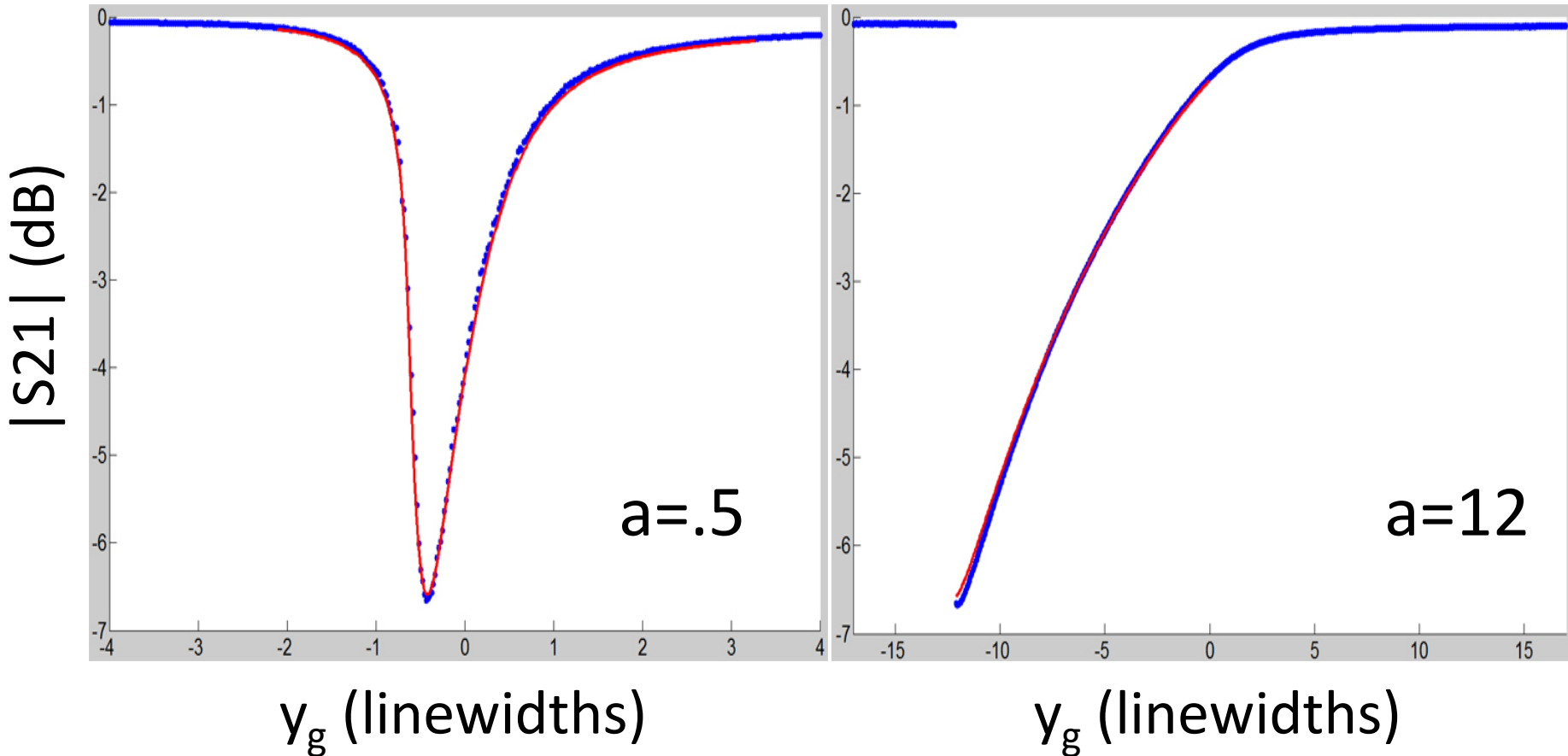
High power:

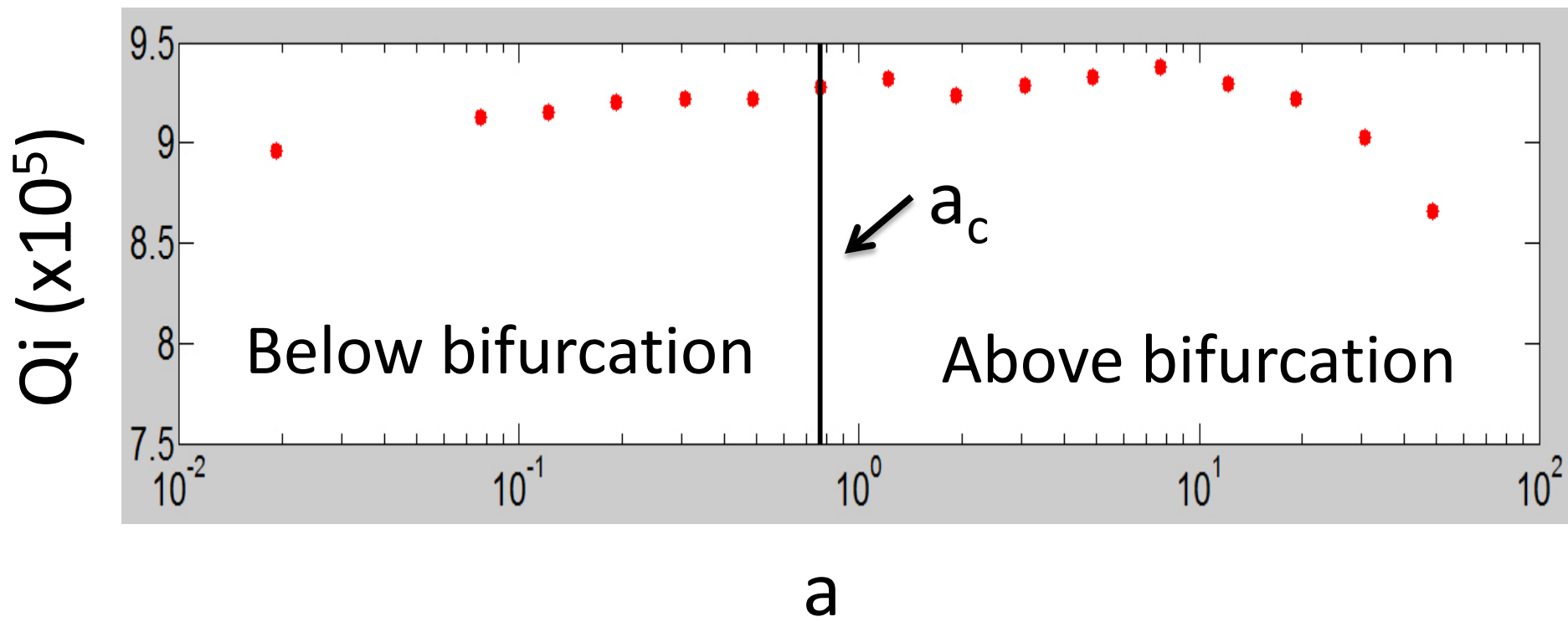
- 1) Fix everything except Q_i by fitting low power curve ($a \sim .01$)
- 2) Fit downward sweeping curves.



Fitting single-tone curves

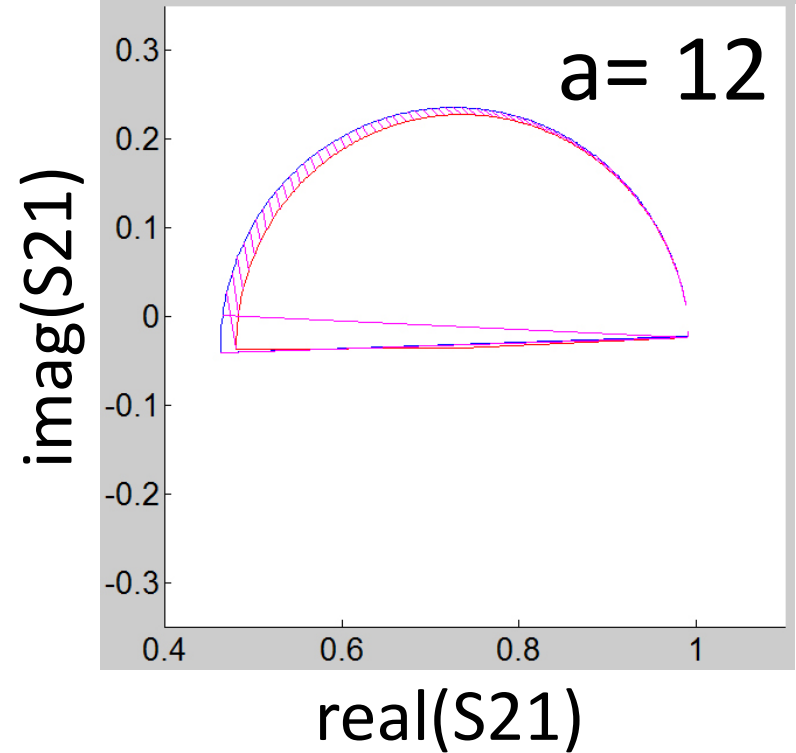
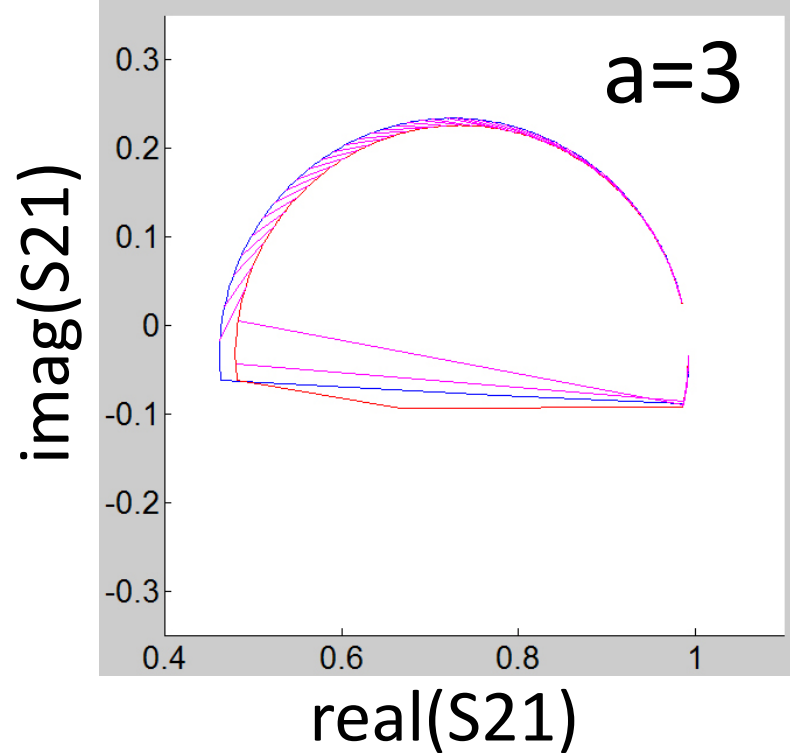
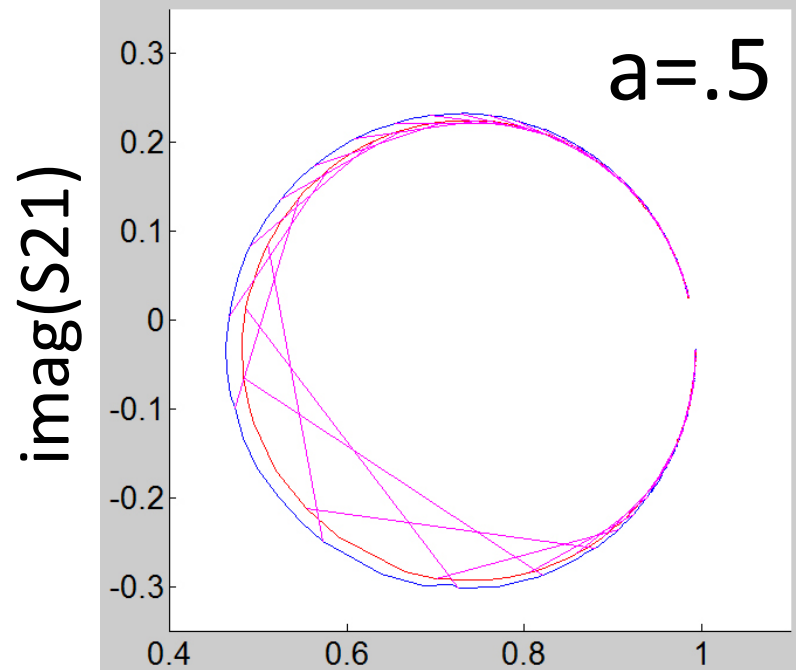
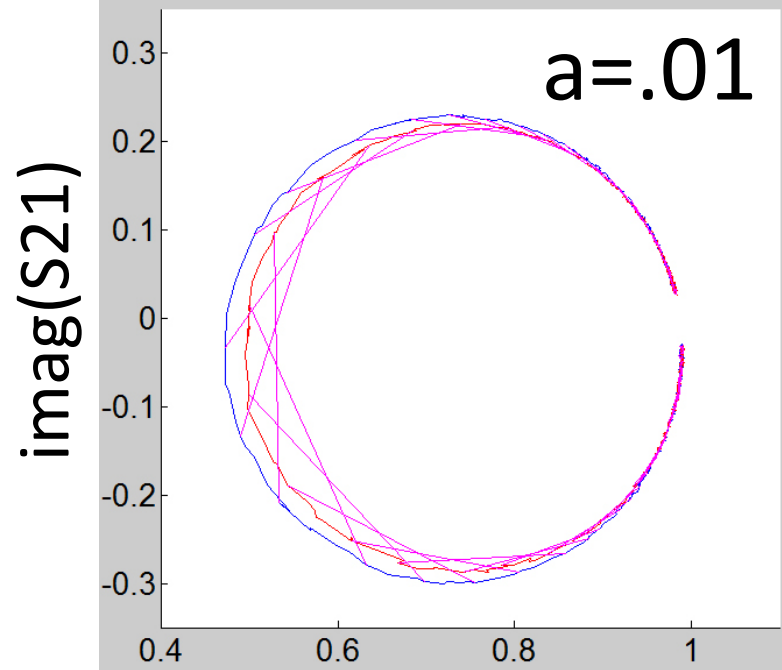
$$t_{21}(f) = ae^{-2\pi jf\tau} \left[1 - \frac{Q_r/Q_c e^{j\phi_0}}{1 + 2j\gamma} \right]$$





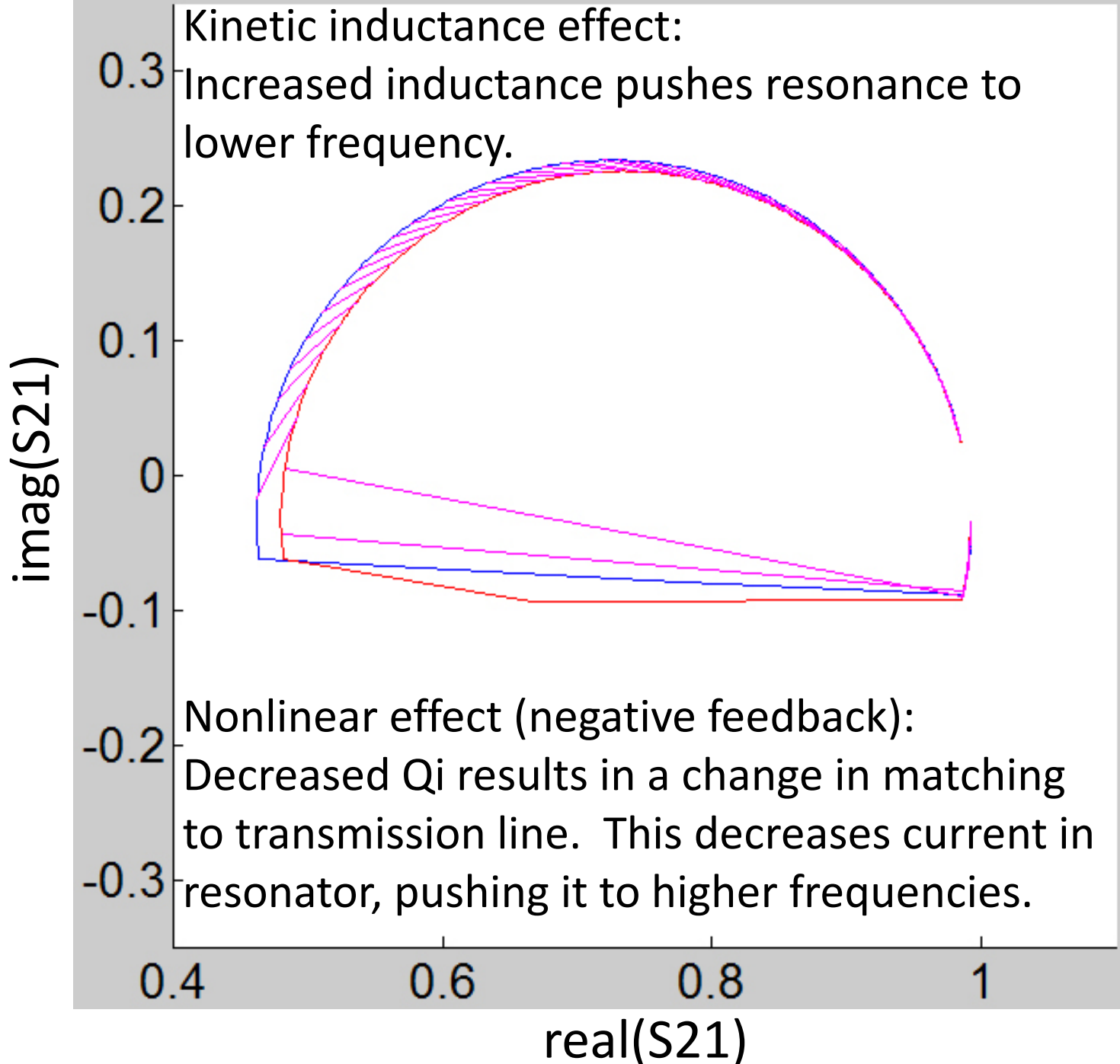
NEP calculations: Signal

Signal: Response to .5 pW while
under 5 pW loading



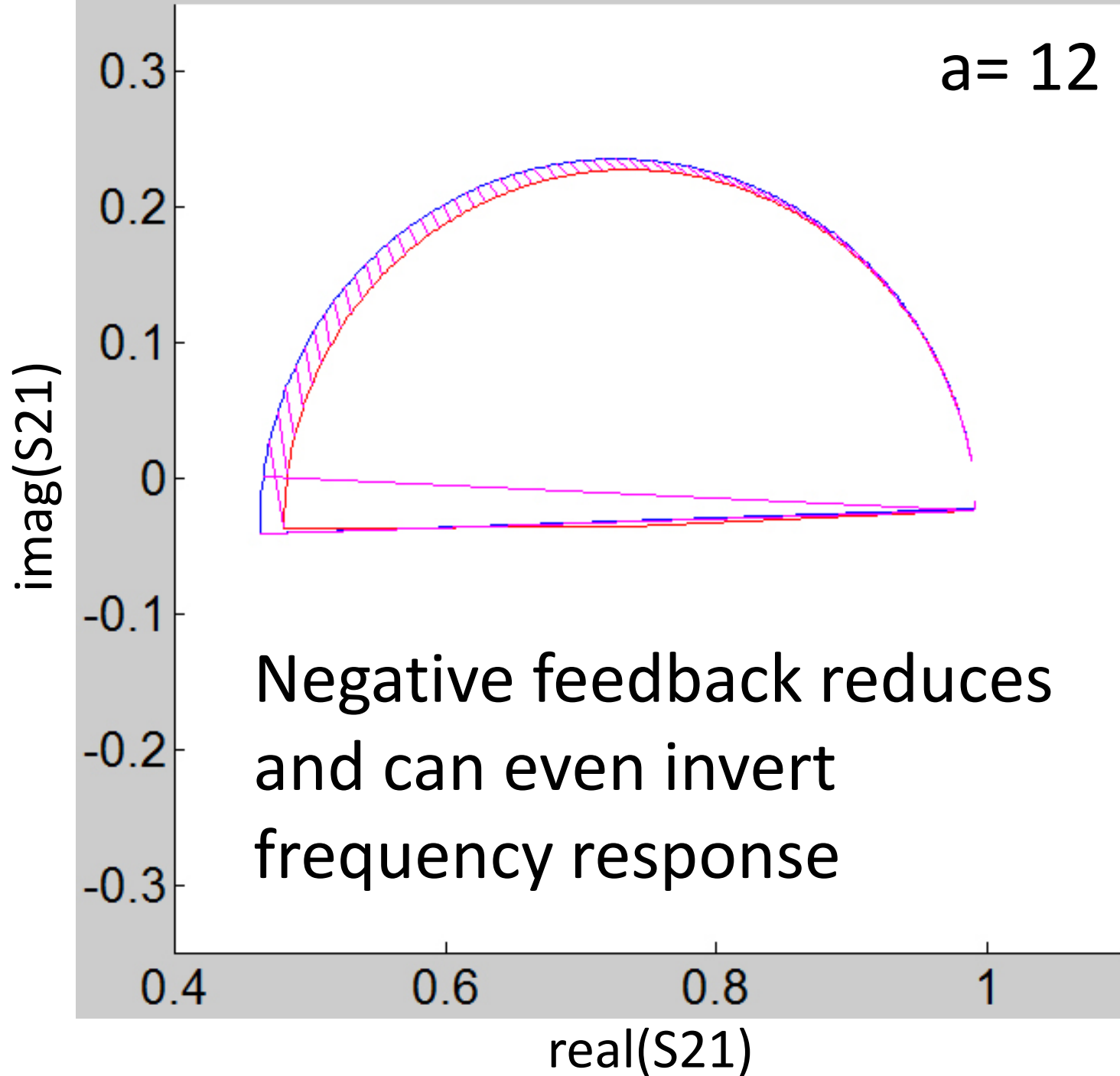
a=3

Kinetic inductance effect:
Increased inductance pushes resonance to lower frequency.



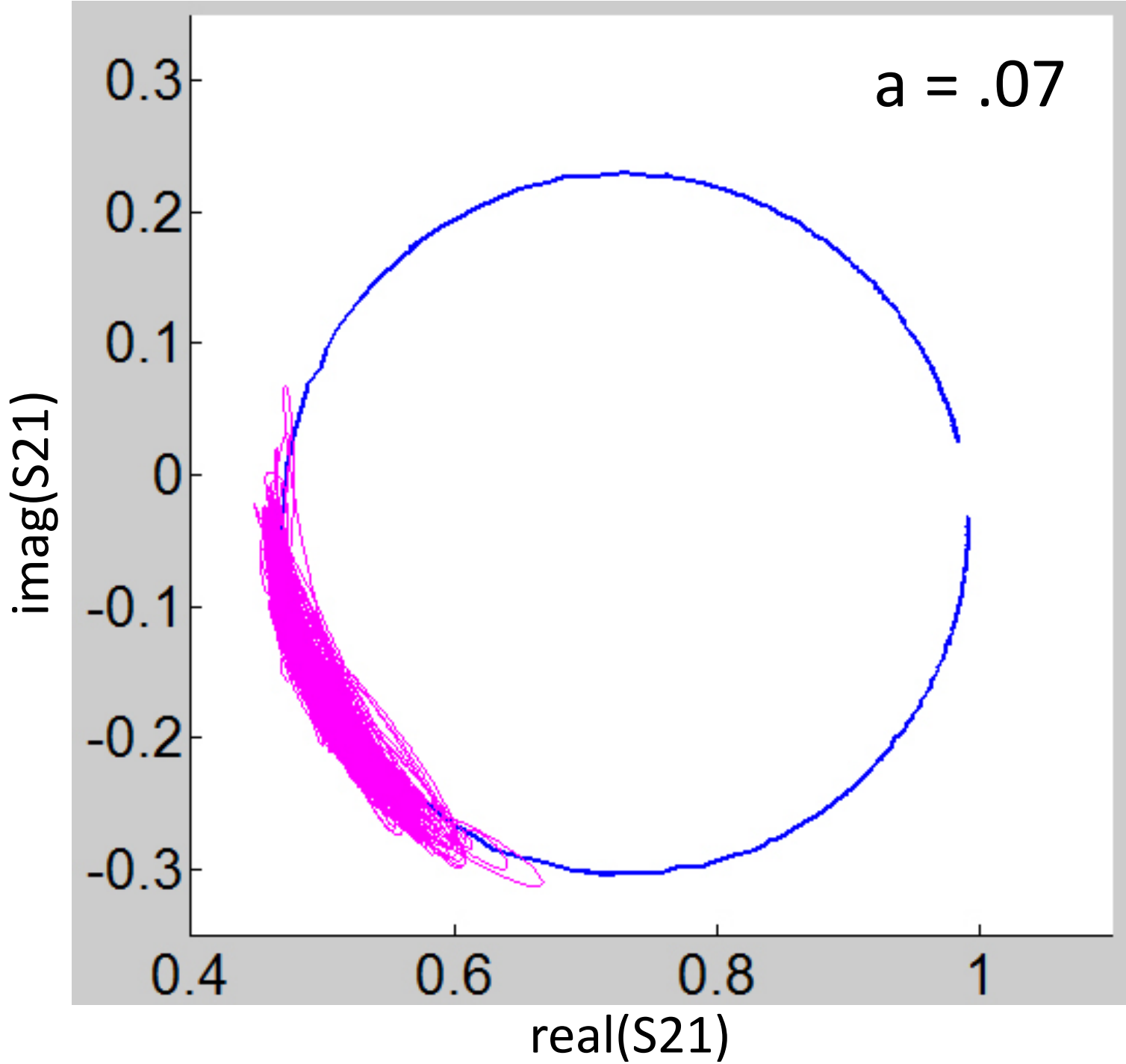
Nonlinear effect (negative feedback):
Decreased Q_i results in a change in matching to transmission line. This decreases current in resonator, pushing it to higher frequencies.

a= 12



Negative feedback reduces
and can even invert
frequency response

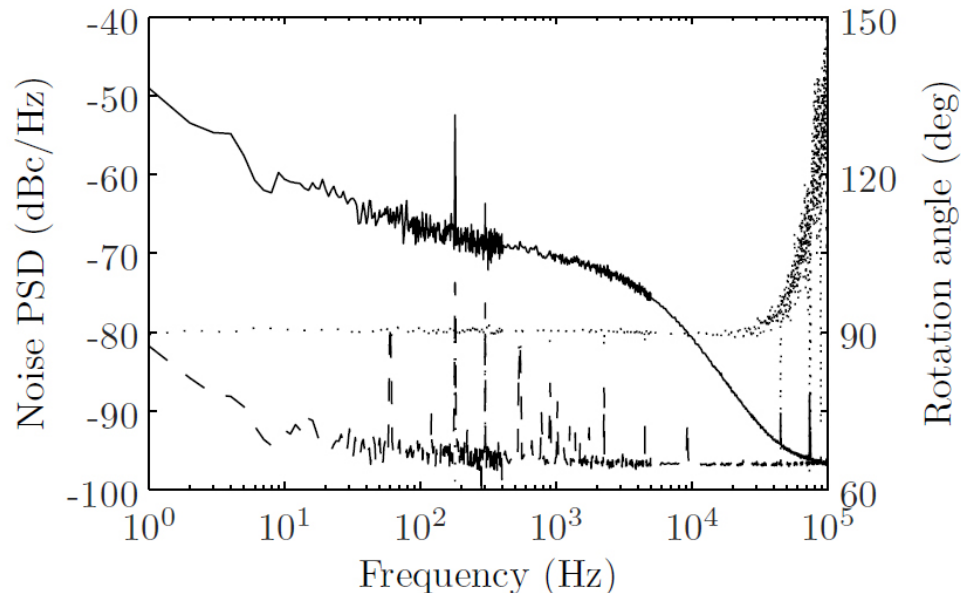
NEP calculations: Noise



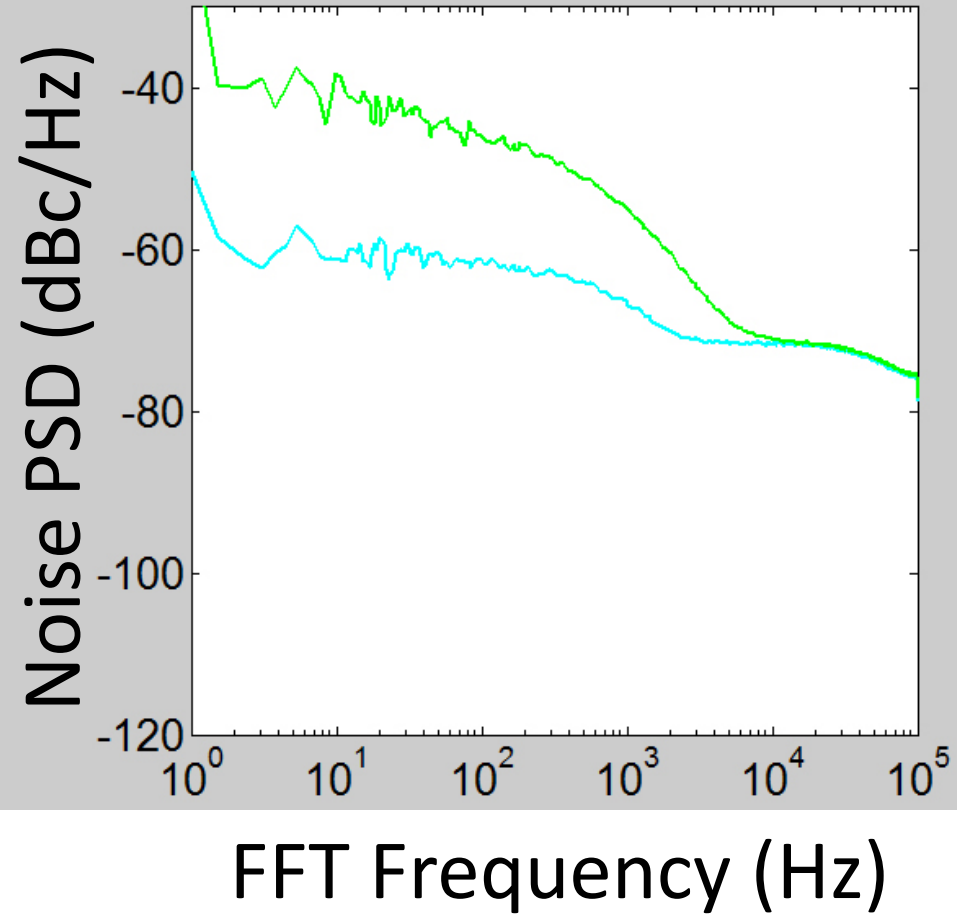
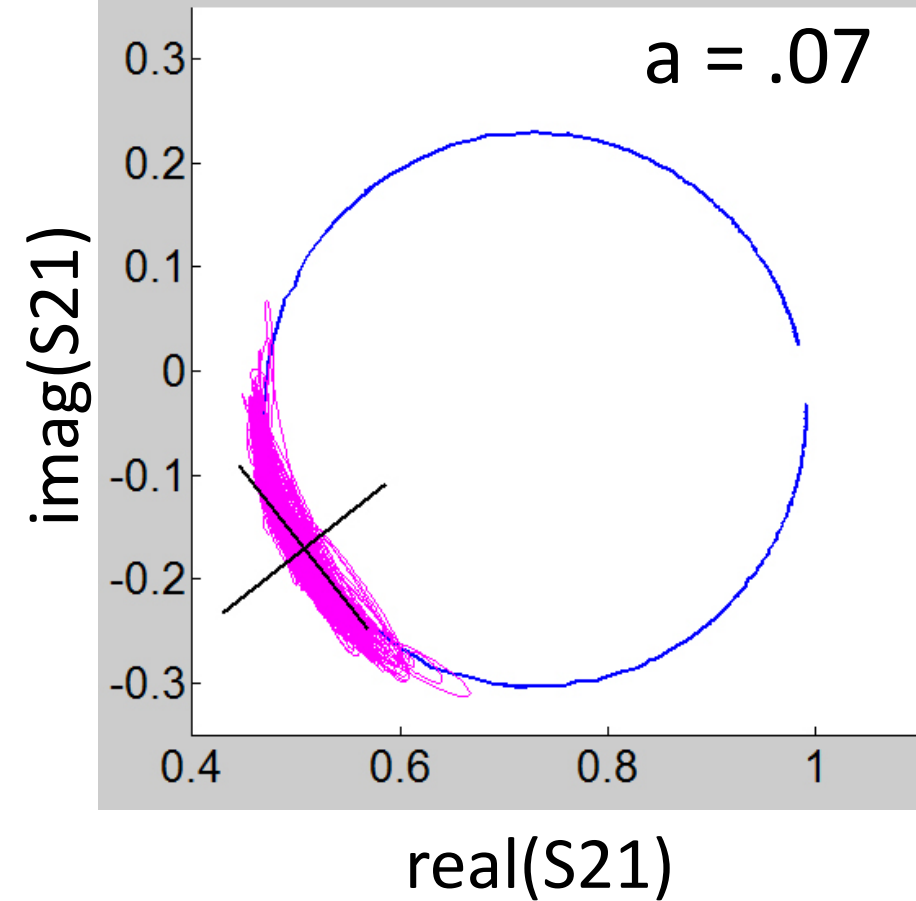
Jiansong Gao Method*: Diagonalize Spectral Density Matrix

$$\langle \delta\xi(\nu)\delta\xi^\dagger(\nu') \rangle = S(\nu)\delta(\nu - \nu'), \quad S(\nu) = \begin{pmatrix} S_{II}(\nu) & S_{IQ}(\nu) \\ S_{IQ}^*(\nu) & S_{QQ}(\nu) \end{pmatrix} \quad (5.2)$$

$$O^T(\nu) \operatorname{Re} S(\nu) O(\nu) = \begin{pmatrix} S_{aa}(\nu) & 0 \\ 0 & S_{bb}(\nu) \end{pmatrix} \quad (5.3)$$



Diagonalized noise spectra



At low powers, noise has significant curvature

Transform to remove curvature

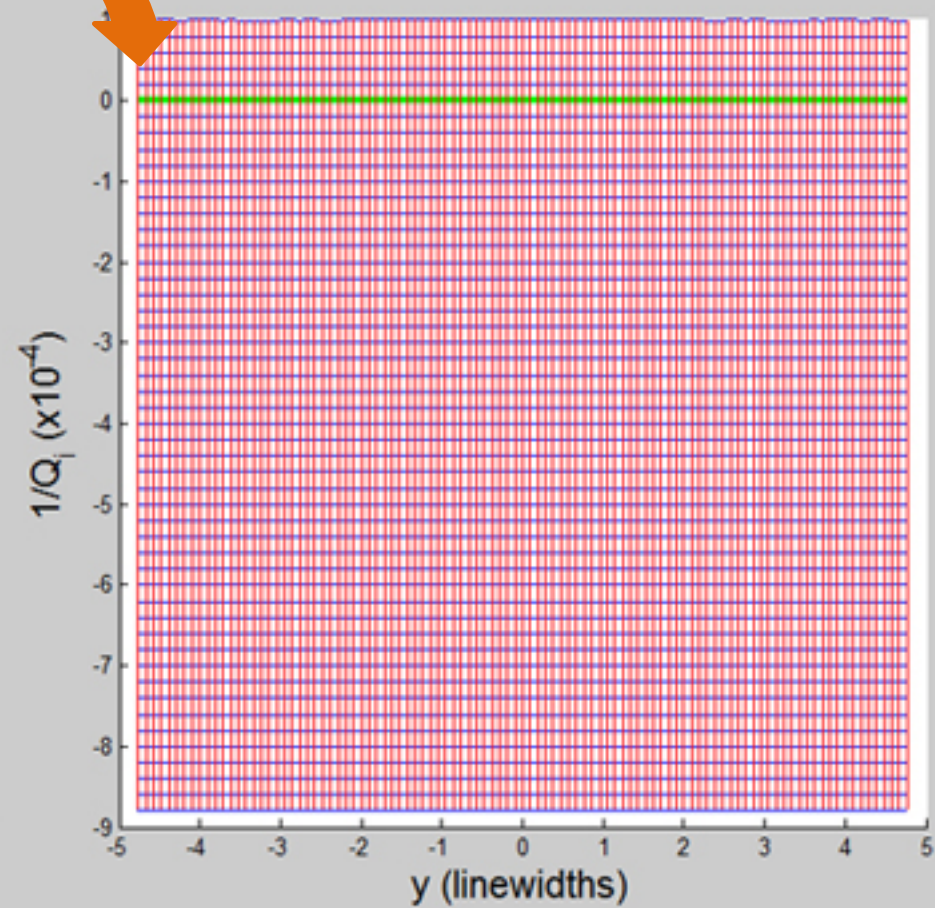
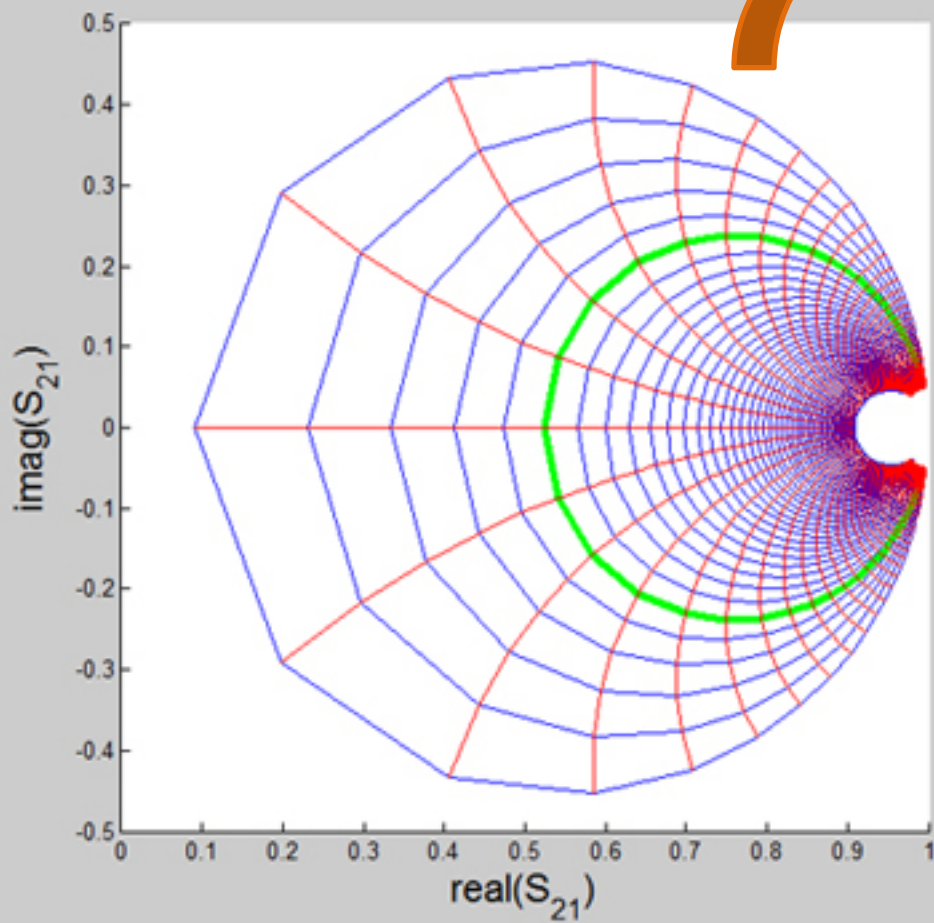
$$S_{21} = 1 - (Q_r/Q_c)[1/(1+2jQ_r x)]$$

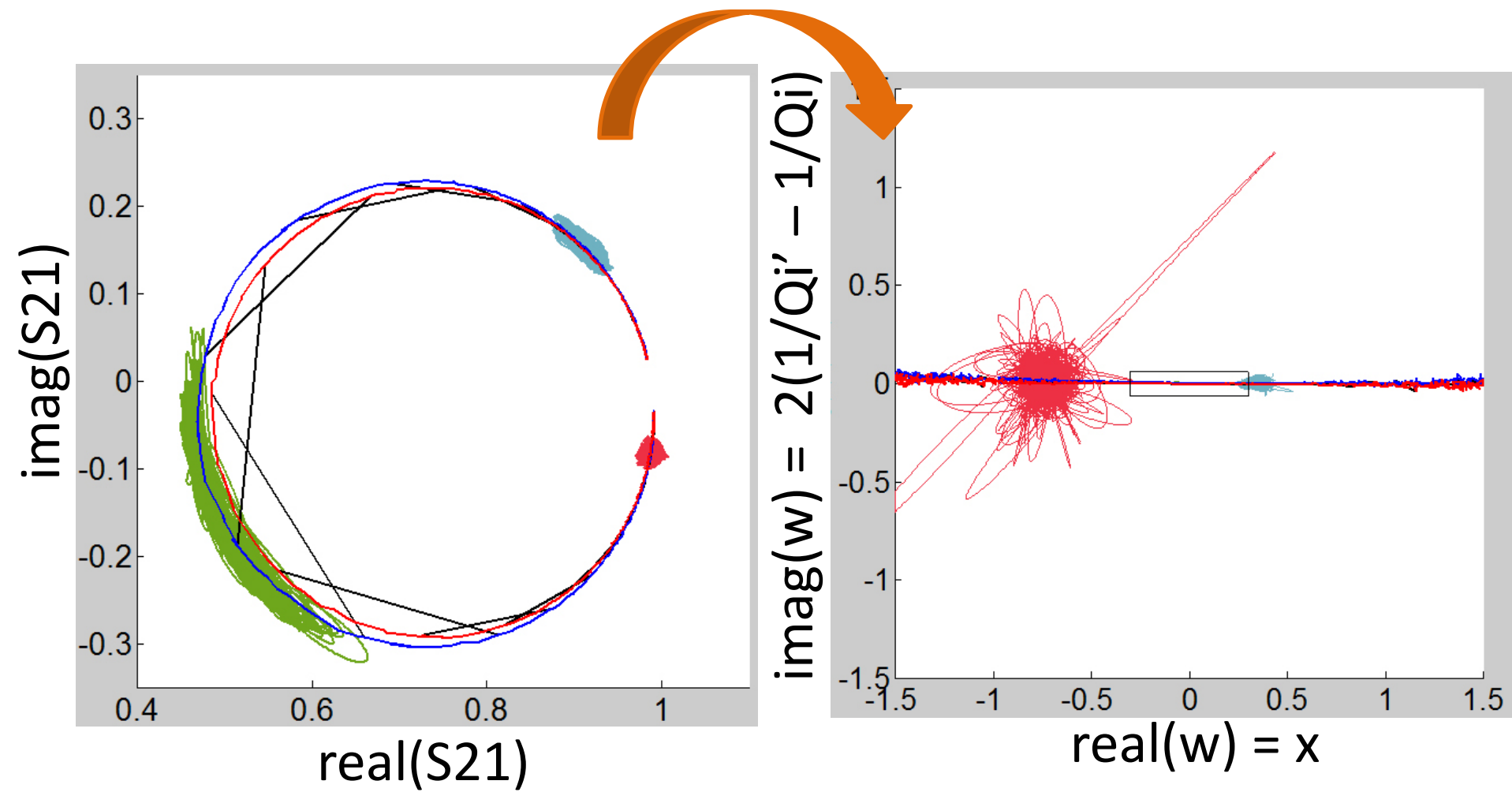
$$S_{21} = 1 - (Q_r/Q_c)[1/(1+2jQ_r \mathbf{w})]$$

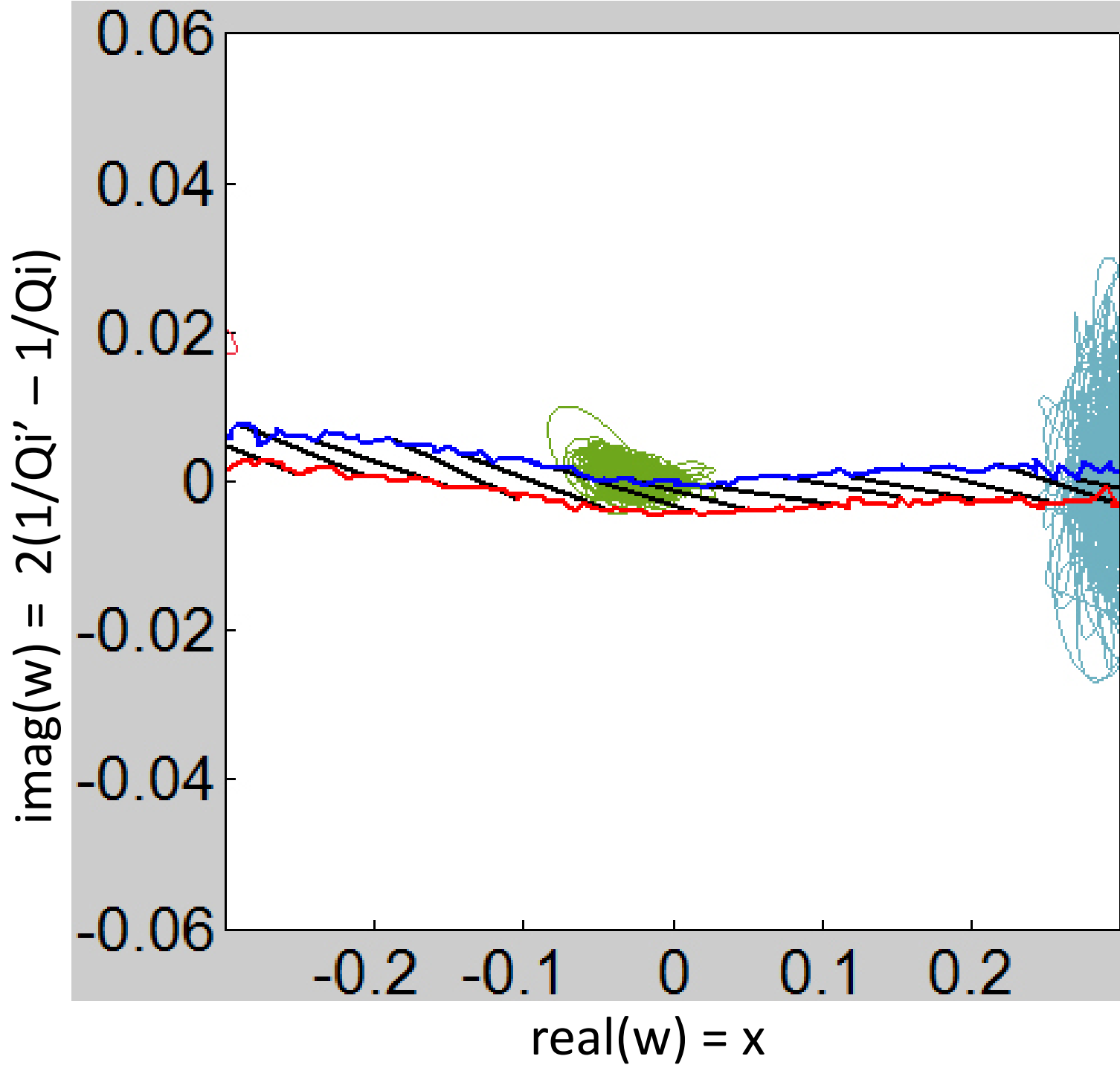
$$\mathbf{w} = 1/(2jQ_c) (1/(1 - S_{21}) - 1/(2jQ_r)); \text{ fixed } Q_i'$$

$$\text{real}(\mathbf{w}) = x; \text{ imag}(\mathbf{w}) = 2(1/Q_i' - 1/Q_i)$$

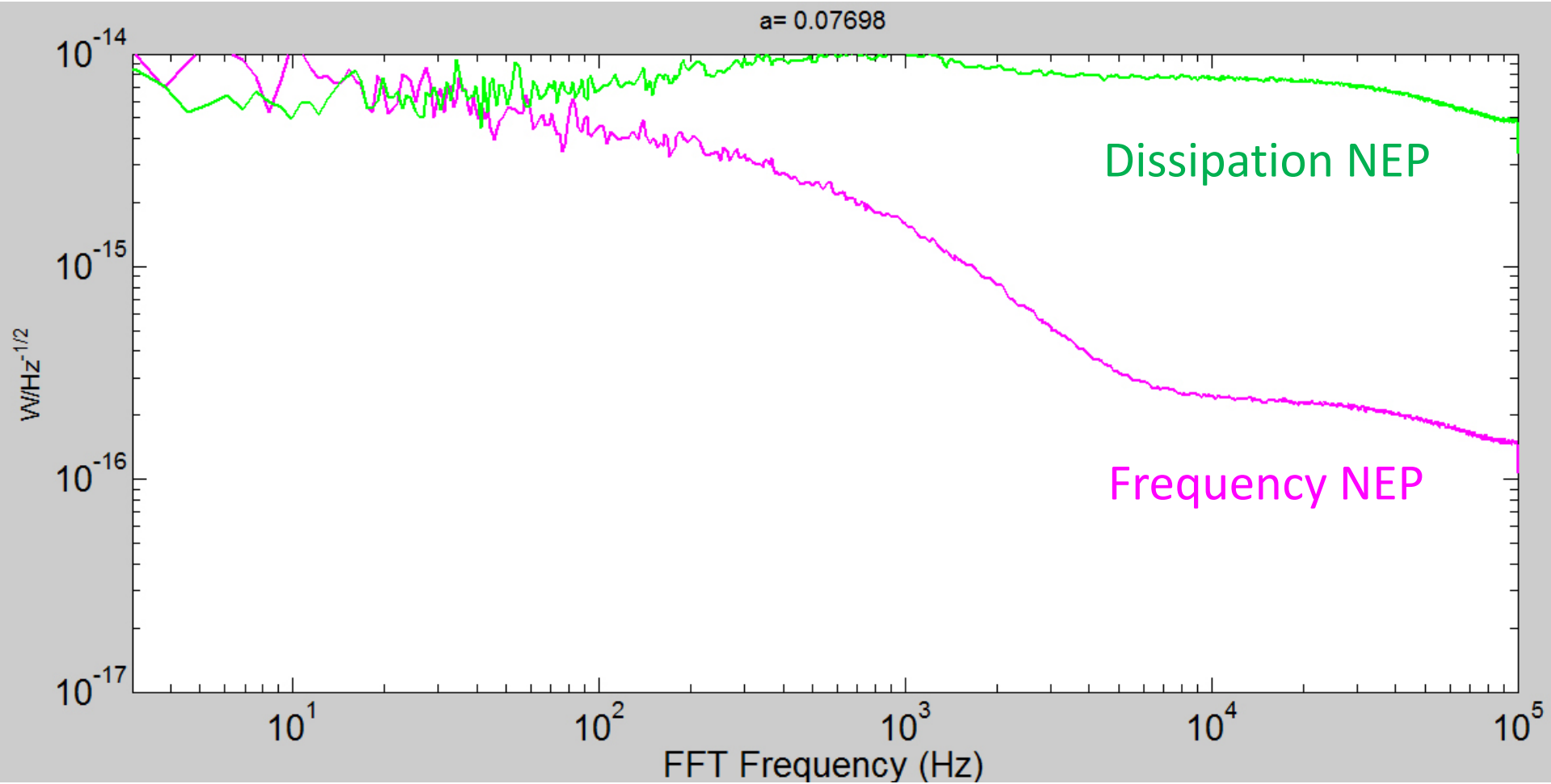
Transform





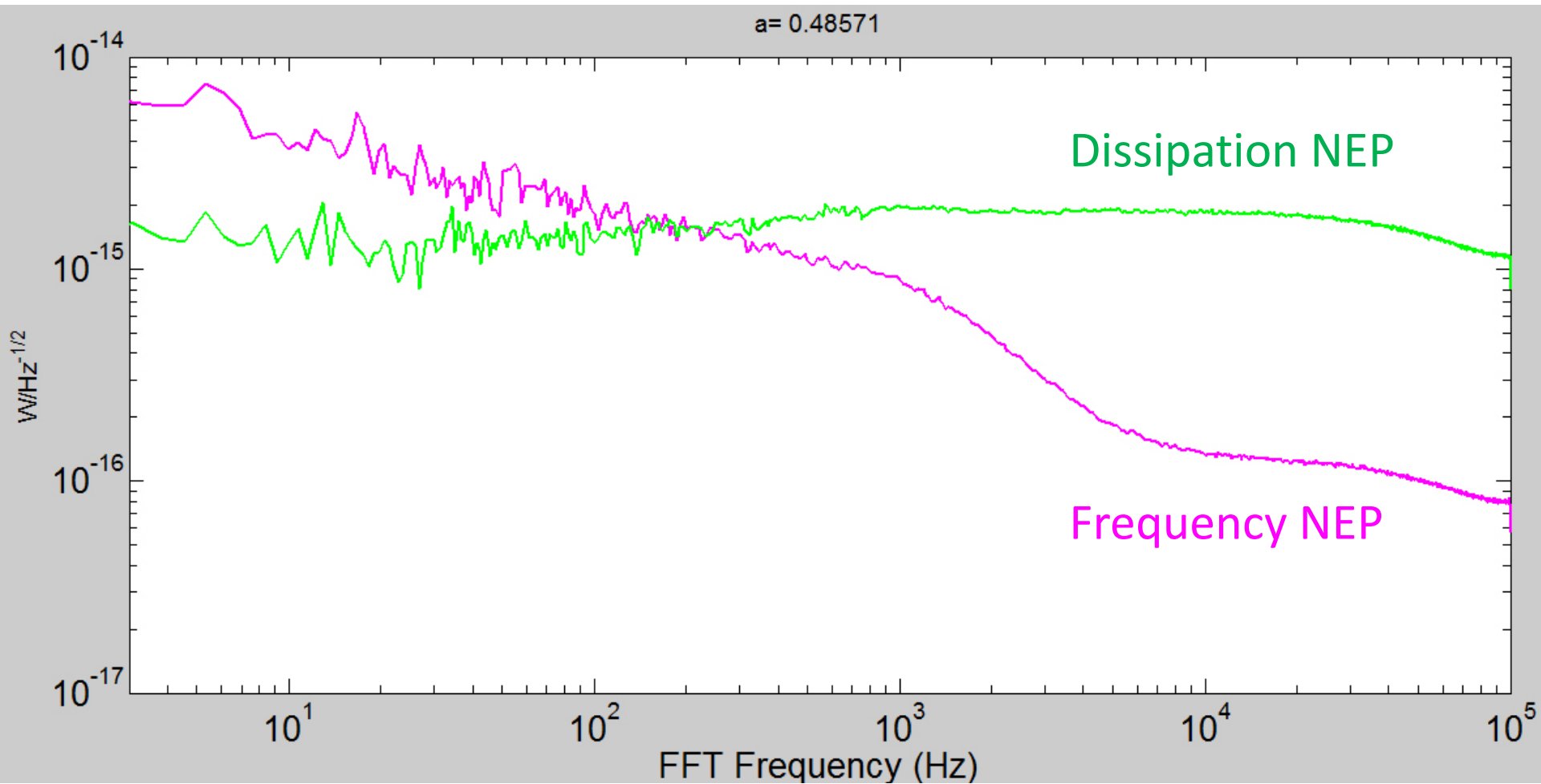


a = .08



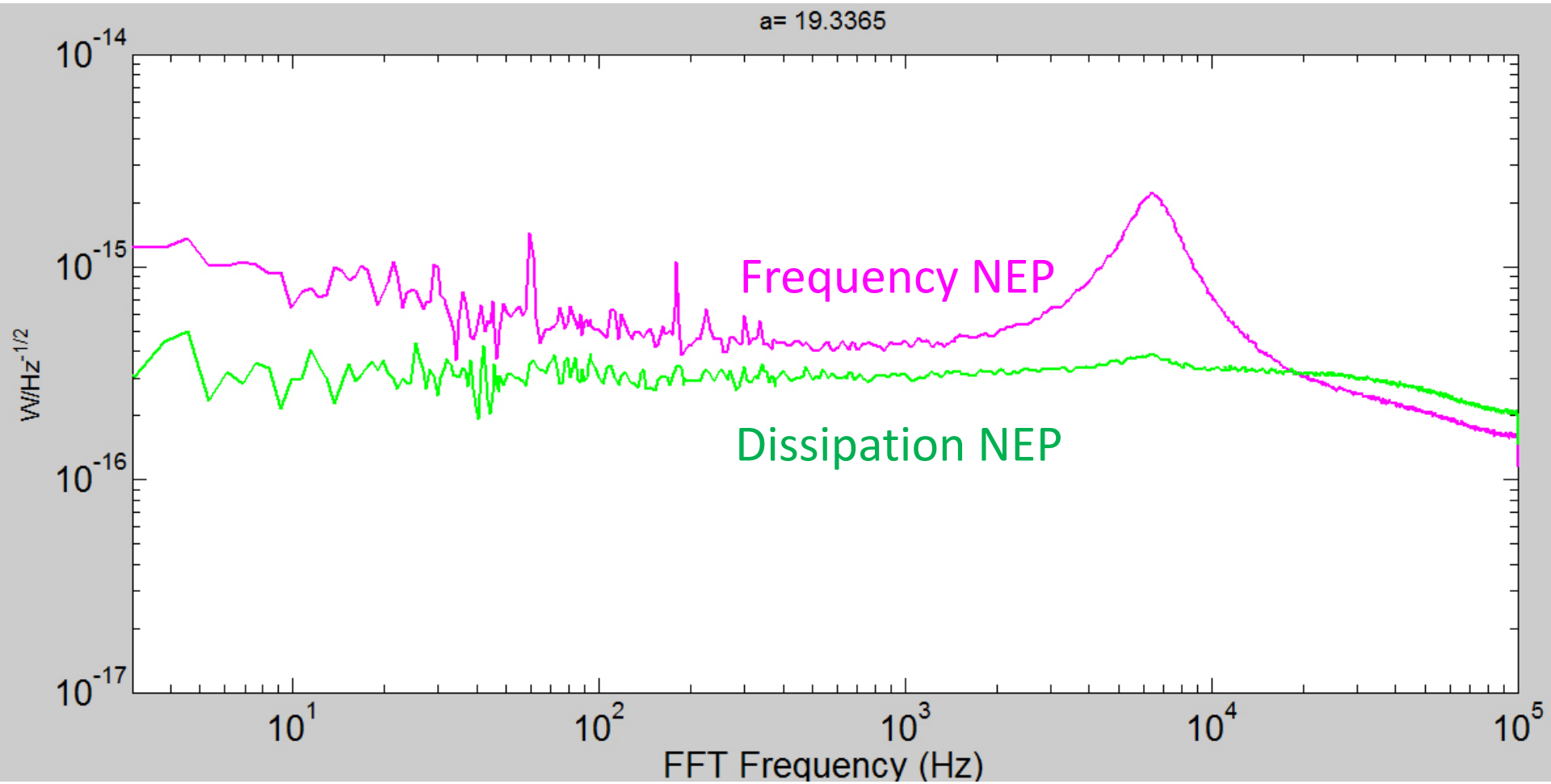
$$\text{NEP} = 6 \times 10^{-15} \text{ W/Hz}^{1/2}$$

a = .5



$$\text{NEP} = 1 \times 10^{-15} \text{ W/Hz}^{1/2}$$

a = 19



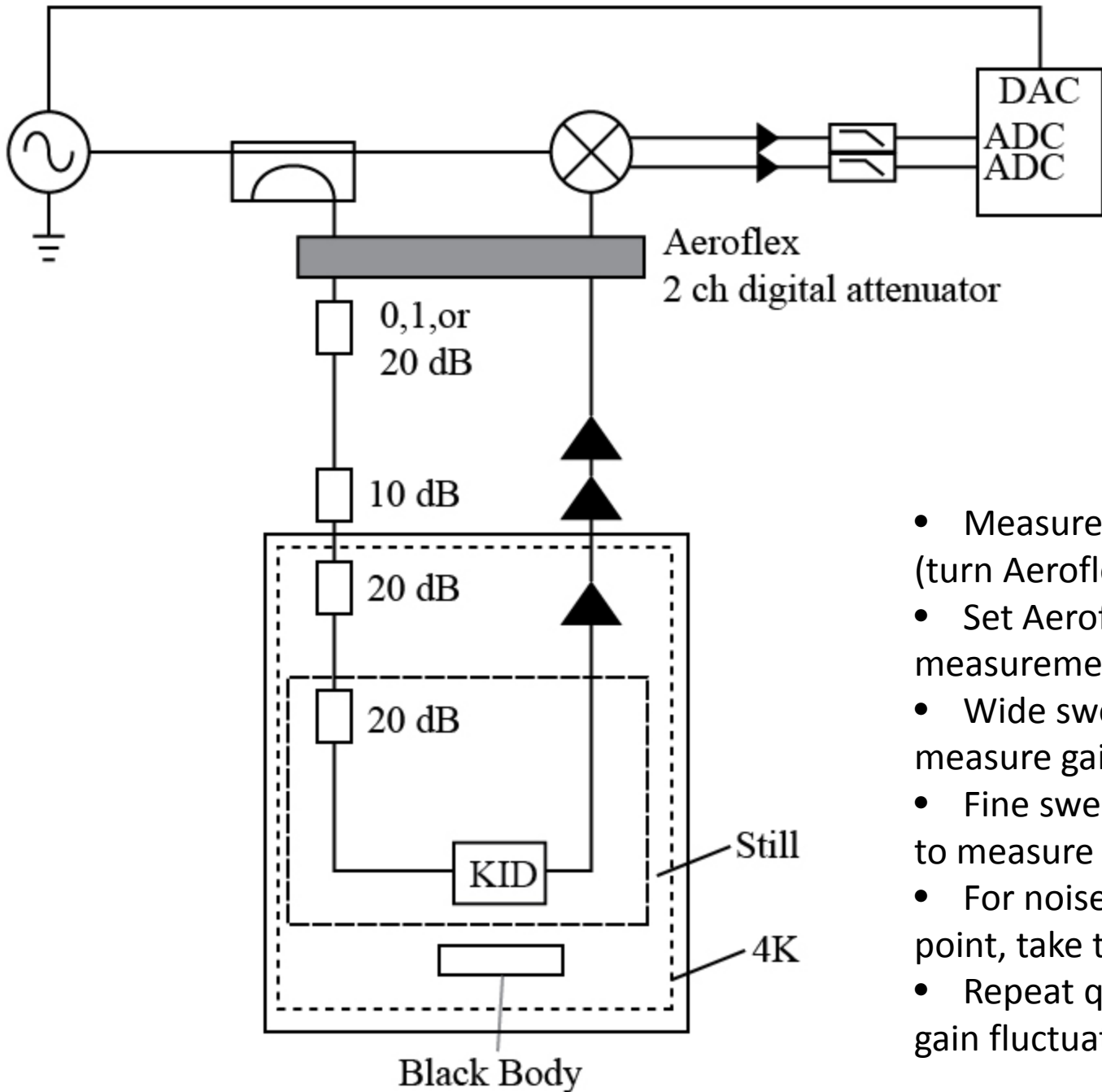
$$\text{NEP} = 3 \times 10^{-16} \text{ W/Hz}^{1/2}$$

Is this useful?

- Possible to access higher powers by downward frequency sweeping.
- Results in reduced noise, improved NEP.
- In order to operate above bifurcation, it is necessary to downward frequency sweep to operating point. -> Significantly increases electronic complexity, but feasible.
- Reduced dynamic range?
 - > avoid jumps
 - > may be difficult to operate with large sky noise

Thank you!

- Peter Day
- Byeong Ho Eom
- Rick Leduc
- Chris McKenney
- Omid Noroozian
- Jonas Zmuidzinas



- Measure DC offset (turn Aeroflex to max attenuation)
- Set Aeroflex to appropriate measurement value
- Wide sweep the resonance to measure gain, cable delay
- Fine sweep the resonance to measure signal
- For noise, fine sweep to noise point, take time trace
- Repeat quickly (minimize gain fluctuations, etc)